# Series of Sobolev Inequalities with Remainder Terms 

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## Abstract

The Series of Sobolev inequality in $\mathbb{R}^{3+\epsilon}, \epsilon \geq 0$, asserts that $\left\|\sum \nabla f_{j}\right\|_{2}^{2} \geq \mathbb{S}_{3+\epsilon}\left(\sum\left\|f_{j}\right\|_{\left.\frac{2(3+\epsilon)}{2}\right) \text {, with } \mathbb{S}_{3+\epsilon} \text {, being } 1+\epsilon}\right.$ the sharp constant. This paper is concerned, with functions restricted to bounded domains $\Omega \subset \mathbb{R}^{3+\epsilon}$. Following H. Brezis, E. Lieb [13] two kinds of inequalities are established: (i) If $f_{j}=0$ on $\partial \Omega$, then $\left\|\Sigma \nabla f_{j}\right\|_{2}^{2} \geq$ $\mathbb{S}_{3+\epsilon}\left(\sum\left\|f_{j}\right\|_{\left.\frac{2(3+\epsilon)}{2}\right)}^{1+\epsilon}\right)+C(\Omega)\left(\sum\left\|f_{j}\right\|_{\frac{(3+\epsilon)}{1+\epsilon}, w}^{2}\right)$ and $\Sigma\left\|\nabla f_{j}\right\|_{2}^{2} \geq \mathbb{S}_{3+\epsilon}\left(\sum\left\|f_{j}\right\|_{2^{*}}^{2}\right)+D(\Omega)\left(\sum\left\|f_{j}\right\|_{\frac{3+\epsilon}{2+\epsilon} w}^{2}\right)$. ii) If $f_{j} \neq 0$ on $\partial \Omega$, then $\sum\left\|\nabla f_{j}\right\|_{2}+C(\Omega)\left(\sum\left\|f_{j}\right\|_{\frac{3+\epsilon}{2+\epsilon}, \Omega}\right) \geq S_{3+\epsilon}^{1 / 2}\left(\sum\left\|f_{j}\right\|_{\frac{2(3+\epsilon)}{1+\epsilon}}\right)$ with $\epsilon^{2}+a \epsilon+5=0$. Some further results and open problems in this area are also presented.

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## I. Introduction

The usual Series of Sobolev Sobolev inequality in $\mathbb{R}^{3+\epsilon}, \epsilon \geq 0$, for the $\mathrm{L}^{2}$ norm of the gradient is

$$
\begin{equation*}
\left\|\sum \nabla f_{j}\right\|_{2}^{2} \geq S_{3+\epsilon}\left(\sum\left\|f_{j}\right\|_{\frac{2(3+\epsilon)}{2}}^{2+\epsilon}\right), \tag{1.1}
\end{equation*}
$$

for all functions $\mathrm{f}_{\mathrm{j}}$ with $\sum \nabla \mathrm{f}_{\mathrm{j}} \in \mathrm{L}^{2}$ and with $\mathrm{f}_{\mathrm{j}}$ vanishing at infinity in the weak sense that means $\left\{\mathrm{x}\left|\left|\mathrm{f}_{\mathrm{j}}(\mathrm{x})\right|>\right.\right.$ $\mathrm{a}<\infty$ for all $\mathrm{a}>O$ (see [12]). The sharp constant $S 3+\epsilon$, is known to be

$$
\begin{equation*}
s_{3+\epsilon}=\pi(3+\epsilon)(1-2)[\Gamma((3+\epsilon) / 2) / \Gamma(3+\epsilon)]^{2 / 3+\epsilon} . \tag{1.2}
\end{equation*}
$$

The constant $S_{3+\epsilon}$, is achieved in (1.1) if and only if

$$
\begin{equation*}
\mathrm{f}_{\mathrm{j}}(\mathrm{x})=\mathrm{a}\left[\varepsilon^{2}+|\epsilon|^{2}\right]^{-(1+\epsilon) / 2} \tag{1.3}
\end{equation*}
$$

for some $\mathrm{a} \in \mathbb{C}, \varepsilon \neq 0$ and $(\mathrm{x}+\epsilon) \in \mathbb{R}^{3+\epsilon}[1,2,6,7,9,11]$.
We consider appropriate modifications of (1.1) when $\mathbb{R}^{3+\epsilon}$ is replaced by a bounded domain $\Omega \subset \mathbb{R}^{3+\epsilon}$. There are two main problems (See [13]):
Problem A. If $\sum \mathrm{f}_{\mathrm{j}}=0$ on $\partial \Omega$, then (1.1) still holds (with $L^{\frac{(3+\epsilon)}{1+\epsilon}}$ norms in $\Omega$, of course), since $\mathrm{f}_{\mathrm{j}}$ can be extended to be zero outside of $\Omega$. In this case (1.1) becomes a strict inequality when $\sum \mathrm{f}_{\mathrm{j}} \neq 0$ (in view of (1.3). However, $\mathrm{S}_{3+\epsilon}$, is still the sharp constant in (1.1) (since $\sum\left\|\nabla \mathrm{f}_{\mathrm{j}}\right\|_{2} /\left\|\mathrm{f}_{\mathrm{j}}\right\|_{\frac{2(3+\epsilon)}{1+\epsilon}}$ is scale invariant). Our goal, in this case, is to give a lower bound to the difference of the two sides in (1.1) for $\mathrm{f}_{\mathrm{j}} \in \mathrm{H}_{0}^{1}(\Omega)$. In Section II we shall prove the following inequalities (1.4) and (1.6):

$$
\begin{equation*}
\sum\left\|\nabla f_{j}\right\|_{2}^{2} \geq S_{3+\epsilon}\left(\sum\left\|f_{j}\right\|_{2}^{2}\right)+C(\Omega)\left(\sum\left\|f_{j}\right\|_{\frac{3+\epsilon}{1+\epsilon}, \mathrm{w}}^{2}\right) \tag{1.4}
\end{equation*}
$$

Where $\mathrm{C}(\Omega)$ depends on $\Omega$ and $3+\epsilon, \frac{3+\epsilon}{1+\epsilon}$, and w denotes the weak $L^{\frac{3+\epsilon}{1+\epsilon}}$ norm defined by

$$
\sum\left\|\mathrm{f}_{\mathrm{j}}\right\|_{\frac{3+\epsilon}{1+\epsilon^{W}}}=\sup _{\mathrm{A}}|\mathrm{~A}|^{-1 /\left(\frac{3+\epsilon}{1+\epsilon}\right)^{\prime}} \int_{\mathrm{A}} \sum\left|\mathrm{f}_{\mathrm{j}}(\mathrm{x})\right| \mathrm{dx},
$$

With A being a set of finite measure $|\mathrm{A}|$.
The inequality (1.4) was motivated by the weaker inequality in [3],

$$
\begin{equation*}
\sum\left\|\nabla f_{j}\right\|_{2}^{2} \geq S_{3+\epsilon}\left(\sum\left\|f_{j}\right\|_{\frac{2(3+\epsilon)}{1+\epsilon}}^{2}\right)+C_{\frac{3+\epsilon}{1+\epsilon}}(\Omega)\left(\sum\left\|f_{j}\right\|_{\frac{3+\epsilon}{1+\epsilon}}^{2}\right), \tag{1.5}
\end{equation*}
$$

which holds for all $\frac{3+\epsilon}{1+\epsilon}$ (with $\mathrm{C}_{\frac{3+\epsilon}{}}^{1+\epsilon}(\Omega) \rightarrow 0$ as $\frac{2(3+\epsilon)}{1+\epsilon}$ ). The proof of (1.5) in [3] was very indirect compared to the proof of (1.4) given here. Inequality (1.4) is best possible in the sense that (1.5) cannot hold with $\frac{3+\epsilon}{1+\epsilon}$; this can be shown by taking the $f_{j}$ in (1.3), applying a cutoff function to make $f_{j}$ vanish on the boundary, and then expanding the integrals (as in [3]) near $\varepsilon=0$.

An inequality stronger than (1.4), and involving the gradient norm is

$$
\begin{equation*}
\left\|\sum \nabla f_{j}\right\|_{2}^{2} \geq S_{3+\epsilon}\left(\sum\left\|f_{j}\right\|_{\frac{2(3+\epsilon)}{1+\epsilon}}^{2}\right)+D(\Omega)\left(\sum\left\|\nabla f_{j}\right\|_{\frac{3+\epsilon}{2+\epsilon}, \mathrm{w}}^{2}\right), \tag{1.6}
\end{equation*}
$$

with $\frac{3+\epsilon}{2+\epsilon}$. (The reason that (1.6) is stronger than (1.4) is that the Sobolev inequality has an extension to the weak norms, by Young's inequalities in weak $L^{\frac{3+\epsilon}{1+\epsilon}}$ spaces).

Among the open questions concerning (1.4)-(1.6) are the following:
(a) What are the sharp constants in (1.4)-(1.6)? Are they achieved? Except in one case, they are not known, even for a ball. If $\epsilon=0, \Omega$ is a ball of radius R and $\epsilon=2$ in (1.6), then $\mathrm{C}_{2}(\Omega)=\pi^{2} /\left(4 \mathrm{R}^{2}\right)$; however, this constant is not achieved [3].
(b) What can replace the right side of (1.4)-(1.6) when $\Omega$ is unbounded, e.g., a half-space?
(c) Is there a natural way to bound $\sum\left\|\nabla \mathrm{f}_{\mathrm{j}}\right\|_{2}^{2}-\mathrm{S}_{3+\epsilon}\left(\sum\left\|\mathrm{f}_{\mathrm{j}}\right\|_{\left.\frac{2(3+\epsilon)}{2}\right)}^{1+\epsilon}\right)$ from below in terms of the "distance" of $f_{j}$ from the set of optimal functions (1.3)?
Problem B. If $\sum \mathrm{f}_{\mathrm{j}} \neq 0$ on $\partial \Omega$, then (1.1) does not hold in $\Omega$ (simply take $\sum \mathrm{f}_{\mathrm{j}}=1$ in $\Omega$ ). Let us assume now that $\Omega$ is not only bounded but that $\partial \Omega$ (the boundary of $\Omega$ ) has enough smoothness. Then (1.1) might be expected to hold if suitable boundary integrals are added to the left side. In Section III we shall prove that for $\sum \mathrm{f}_{\mathrm{j}}=$ constant $\equiv \sum \mathrm{f}_{\mathrm{j}}(\partial \Omega)$ on $\partial \Omega$

$$
\begin{equation*}
\left\|\sum \nabla \mathrm{f}_{\mathrm{j}}\right\|_{2}^{2}+\mathrm{E}(\Omega)\left|\sum \mathrm{f}_{\mathrm{j}}(\partial \Omega)\right|^{2} \geq \mathrm{S}_{3+\epsilon}\left(\sum \mid \mathrm{f}_{\mathrm{j}} \|_{\left.\frac{2(3+\epsilon)}{2}\right)}^{2+\epsilon} .\right. \tag{1.7}
\end{equation*}
$$

On the other hand, if $\mathrm{f}_{\mathrm{j}}$ is not constant on $\partial \Omega$, then the following two inequalities hold.

$$
\begin{align*}
\left\|\sum \nabla f_{j}\right\|_{2}^{2}+F(\Omega)\left(\left\|\sum f_{j}\right\|_{H^{1 / 2}(\partial \Omega)}^{2}\right. & \geq S_{3+\epsilon}\left(\sum\left\|f_{j}\right\|_{\left.\frac{2(3+\epsilon)}{2}\right)}^{2}\right),  \tag{1.8}\\
\left\|\sum \nabla f_{j}\right\|_{2}+G(\Omega)\left(\left\|\sum f_{j}\right\|_{\frac{3+\epsilon}{2+\epsilon} \partial \Omega}\right) & \geq S_{3+\epsilon}^{1 / 2}\left(\sum\left\|f_{j}\right\|_{\frac{2(3+\epsilon)}{1+\epsilon}}\right) \tag{1.9}
\end{align*}
$$

with $\epsilon^{2}+4 \epsilon+5=0$, which is sharp. (Note the absence of the exponent 2 in (1.9)).
In addition to the obvious analogues of questions (a)-(c) for Problem B, one can also ask whether (1.9) can be improved to

$$
\begin{equation*}
\left\|\sum \nabla \mathrm{f}_{\mathrm{j}}\right\|_{2}^{2}+\mathrm{H}(\Omega)\left(\left\|\sum \mathrm{f}_{\mathrm{j}}\right\|_{\frac{3+\epsilon}{2+\epsilon} \partial \Omega}^{2}\right) \geq \mathrm{S}_{3+\epsilon}\left(\sum\left\|\mathrm{f}_{\mathrm{j}}\right\|_{\left.\frac{2(3+\epsilon)}{1+\epsilon}\right)}^{2}\right) . \tag{1.10}
\end{equation*}
$$

We do not know.
If $\Omega$ is a ball of radius $R$, we shall establish that the sharp constant in (1.7) is $\mathrm{E}(\Omega)=\sigma_{3+\epsilon} \mathrm{R}^{1+\epsilon} /(1+\epsilon)$, where $\sigma_{3+\epsilon}$ is the surface area of the ball of unit radius in $\mathbb{R}^{3+\epsilon}$. With this $\mathrm{E}(\Omega)$, (1.7) is a strict inequality. Given this fact, one suspects (in view of the solution to Problem A) that some term could be added to the right side of (1.5). However, such a term cannot be any $L^{\frac{3+\epsilon}{1+\epsilon}}(\Omega)$ norm of $f_{j}$, as will be shown.

To conclude this Introduction, let us mention two' related inequalities. First, if one is willing to replace $S_{3+\epsilon}$, on the right side of (1.10) by the smaller constant $2^{-2 / 3+\epsilon} S_{3+\epsilon}$, then for a ball one can obtain the inequality

$$
\begin{equation*}
\int \sum\left|\nabla \mathrm{f}_{\mathrm{j}}\right|^{2}+\mathrm{I}(\Omega)\left(\sum\left\|\mathrm{f}_{\mathrm{j}}\right\|_{2, \partial \Omega}^{2}\right) \geq \mathrm{S}^{-2 / 3+\epsilon} \mathrm{S}_{3+\epsilon}\left(\sum\left\|\mathrm{f}_{\mathrm{j}}\right\|_{\left.\frac{2(3+\epsilon)}{2}\right)}^{1+\epsilon}\right) \tag{1.11}
\end{equation*}
$$

This is proved in Section (1.1). Inequalities related to (1.11) were derived by Cherrier [4] for general manifolds.
Second, one can consider the doubly weighted Hardy-Littlewood-Sobolev inequality [7,10] which in some sense is the dual of (1.1), namely,

$$
\begin{gather*}
\left.\left|\iint \sum f_{j}(x) f_{j}(x+\epsilon)\right| \epsilon\right|^{-\lambda}|x|^{-\alpha}|x+\epsilon|^{-\alpha} d x d(x+\epsilon) \mid \\
\leq P_{\alpha, \lambda, 3+\epsilon}\left(\sum\left\|f_{j}\right\|_{\frac{3+\epsilon}{1+\epsilon}}^{2}\right), \tag{1.12}
\end{gather*}
$$

with $\left(\frac{3+\epsilon}{1+\epsilon}\right)^{\prime}=23+\epsilon /(\lambda+2 \alpha), 0<\lambda<3+\epsilon, 0 \leq \alpha<3+\epsilon /\left(\frac{3+\epsilon}{1+\epsilon}\right)^{\prime}$. If $f_{j}$ is restricted to have support in a bounded domain $\Omega$ and if $P$ is (by definition) the sharp constant in $\mathbb{R}^{3+\epsilon}$, one should expect to be able to add some additional term to the left side of (1.12). When $\epsilon=2$ this is indeed possible, and the additional term is:

$$
\begin{equation*}
J_{n}|\Omega|^{-\lambda / 3+\epsilon}\left\{\int \sum f_{j}(x)|x|^{-\alpha} d x\right\}^{2} \tag{1.13}
\end{equation*}
$$

This was proved in [5] for $n=3, \lambda=2, \alpha=\frac{1}{2}$, and $\Omega$ being a ball, but the method easily extends (for a ball) to other $3+\epsilon, \lambda$. The result (1.4) further extends to general $\Omega$ (with the same constant $J_{3+\epsilon}$ ) by using the Riesz rearrangement inequality. On the other hand, when $\epsilon \neq 2$, it does not seem to be easy to find the additional term on the left side of (1.12): at least we have not succeeded in doing so. This is an open problem. In particular, in Section III we prove that when $\epsilon=9, \epsilon=0, \lambda=1, \alpha=0$, one cannot even add $\left\|f_{j}\right\|_{1}^{2}$ to the left side of (1.12).

## II. Proof of Inequalities (1.4) and (1.6):

Proof of Inequalities (1.4)(See [13]): By the rearrangement inequality for the $L^{2}$ norm of the gradient we have

$$
\begin{equation*}
\left\|\sum \nabla f_{j}^{*}\right\|_{2} \leq \sum\left\|\nabla f_{j}\right\|_{2} \tag{2.1}
\end{equation*}
$$

(see, e.g., [8]); in addition we have

$$
\begin{align*}
\sum\left\|f_{j}^{*}\right\|_{2^{*}} & =\sum\left\|f_{j}\right\|_{2^{*}} \\
\sum\left\|f_{j}^{*}\right\|_{\frac{3+\epsilon}{1+\epsilon^{\prime}} w} & =\sum\left\|f_{j}\right\|_{\frac{3+\epsilon}{1+\epsilon^{\prime}},}, \tag{2.2}
\end{align*}
$$

Here, $f_{j}^{*}$ denotes the symmetric decreasing rearrangement of the function $f_{j}$ extended to be zero outside $\Omega$. Therefore, it suffices to consider the case in which $\Omega$ is a ball of radius $R$ (chosen to have the same volume as the original domain) and $f_{j}$ is symmetric decreasing.

Let $g_{j} \in(\Omega)$ and define $u_{j}$ to be the solution of

$$
\begin{array}{rlrl}
\Delta u_{j} & =g_{j} & & \text { in } \\
u_{j} & =0 & & \Omega,  \tag{2.3}\\
\text { on } & & \partial \Omega
\end{array}
$$

Let

$$
\phi_{j}(x)= \begin{cases}f_{j}(x)+u_{j}(x)+\left\|u_{j}\right\|_{\infty} & \text { in } \Omega,  \tag{2.4}\\ \left\|u_{j}\right\|_{\infty}(R /|x|)^{n-2} & \text { in } \Omega^{c} .\end{cases}
$$

The Sobolev inequality in all of $\mathbb{R}^{n}$ applied to $\phi_{j}$ yields

$$
\begin{equation*}
\int_{\Omega} \sum\left|\nabla f_{j}+u_{j}\right|^{2}+\left\|u_{j}\right\|_{\infty}^{2} R^{1+\epsilon}(1+\epsilon) \sigma_{3+\epsilon} \geq S_{3+\epsilon}\left(\sum\left\|f_{j}\right\|_{\frac{2(3+\epsilon)}{2+\epsilon}}^{2}\right) \tag{2.5}
\end{equation*}
$$

Since $\sum f_{j} \geq 0$ and $u_{j}+\left\|u_{j}\right\|_{\infty} \geq 0$. Here

$$
\sigma_{3+\epsilon}=2(\pi)^{3+\epsilon / 2} / \Gamma(3+\epsilon / 2)
$$

is the surface area of the unit ball in $\mathbb{R}^{3+\epsilon}$. Therefore, we find

$$
\begin{equation*}
\int \sum\left|\nabla f_{j}\right|^{2}-2 \int \sum f_{j} g_{j}+\int \sum\left|\nabla u_{j}\right|^{2}+k \sum\left\|u_{j}\right\|_{\infty}^{2} \geq \sum\left\|f_{j}\right\|_{\frac{2(3+\epsilon)}{1+\epsilon}}^{2} \tag{2.6}
\end{equation*}
$$

where $k=R^{1+\epsilon}(1+\epsilon) \sigma_{3+\epsilon}$. Replacing $g_{j}$ by $\lambda g_{j}$ and $u_{j}$ by $\lambda u_{j}$ and optimizing with respect to $\lambda$ we obtain

$$
\begin{equation*}
\int \sum\left|\nabla f_{j}\right|^{2} \geq S_{3+\epsilon}\left(\sum\left\|f_{j}\right\|_{\frac{2(3+\epsilon)}{1+\epsilon}}^{2}\right)+\sum\left(\int f_{j} g_{j}\right)^{2} /\left[\int\left|\nabla u_{j}\right|^{2}+k\left\|u_{j}\right\|_{\infty}^{2}\right] \tag{2.7}
\end{equation*}
$$

In inequality (2.7) we can obviously maximize the right side with respect to $g_{j}$. In view of the definition of the weak norm we shall in fact restrict our attention to $g_{j}=1_{A}$, namely, the characteristic function of some set $A$ in $\Omega$. We shall now establish some simple estimates for all the quantities in (2.7) in which $C_{3+\epsilon}$, generically denotes constants depending only on $3+\epsilon$,

$$
\begin{gather*}
\int \sum f_{j} g_{j}=\int_{A} \sum f_{j}  \tag{2.8}\\
\int \sum\left|\nabla u_{j}\right|^{2} \leq C_{3+\epsilon}|A|^{1+2 / 3+\epsilon}  \tag{2.9}\\
\left\|u_{j}\right\|_{\infty} \leq C_{3+\epsilon}|A|^{2 / 3+\epsilon}
\end{gather*}
$$

Indeed we have, by multiplying (2.3) by $u_{j}$ and using Hölder's inequality,

$$
\begin{align*}
& \int \sum\left|\nabla u_{j}\right|^{2}=-\int_{A} \sum u_{j} \leq \sum\left\|u_{j}\right\|_{\frac{2(3+\epsilon)}{1+\epsilon}}|A|^{\frac{5}{2(3+\epsilon)}} \\
& \leq S_{3+\epsilon}^{-1 / 2}\left(\sum\left\|\nabla u_{j}\right\|_{2}|A|^{\frac{5}{2(3+\epsilon)}}\right) \tag{2.11}
\end{align*}
$$

which implies (2.9). Next we have, by comparison with the solution in $\mathbb{R}^{3+\epsilon}$,

$$
\begin{align*}
\left|u_{j}\right| & \leq C_{3+\epsilon}|x|^{-(1+\epsilon)} *\left(1_{A}\right) \\
& \leq C_{3+\epsilon}^{\prime}|A|^{2 / 3+\epsilon} \tag{2.12}
\end{align*}
$$

since the function $|x|^{-(1+\epsilon)}$ belongs to $L_{w}^{-\frac{3+\epsilon}{11+\epsilon}}$. Since $|A| \leq|\Omega|=\sigma_{3+\epsilon} R^{3+\epsilon} / 3+\epsilon$ we obtain

$$
\begin{equation*}
\int \sum\left|\nabla u_{j}\right|^{2}+k \sum\left\|u_{j}\right\|_{\infty}^{2} \leq C_{3+\epsilon}|A|^{4 / 3+\epsilon} R^{1+\epsilon} \tag{2.13}
\end{equation*}
$$

Hence (1.4) has been proved (for all $\Omega$ ) with a constant

$$
\begin{equation*}
C(\Omega)=C_{3+\epsilon}|\Omega|^{-\frac{1+\epsilon}{3+\epsilon}} \tag{2.14}
\end{equation*}
$$

Proof of Inequality (1.6)(See [13]): To a certain extent the previous proof can be imitated except for one important ingredient, namely, the rearrangement technique cannot be used since it is not true that $\left\|\sum \nabla f_{j}\right\|_{\frac{3+\epsilon}{2+\epsilon},} \leq \sum\left\|\nabla f_{j}^{*}\right\|_{\frac{3+\epsilon}{2+\epsilon}, w}$. (However, it is still true that we can replace $f_{j}$ by $\left|f_{j}\right|$ without changing any of the norms in (1.6), and thus we may and still assume that $\sum f_{j} \geq 0$ ). Consequently we have to use a direct approach and the constant $D(\Omega)$ in (1.6) will not depend only on $|\Omega|$; it will in fact depend on the capacity of $\Omega$. It is an open question whether (1.6) holds with $D(\Omega)$ depending only on $|\Omega|$. Our result is that:

$$
\begin{equation*}
D(\Omega)=C_{3+\epsilon} / \operatorname{cap}(\Omega) \tag{2.15}
\end{equation*}
$$

We begin as before with (2.3), but (2.4) is replaced by:

$$
\phi_{j}= \begin{cases}f_{j}+u_{j}+\left\|u_{j}\right\|_{\infty} & \text { in } \Omega  \tag{2.16}\\ \left\|u_{j}\right\|_{\infty} v_{j} & \text { in } \Omega^{c}\end{cases}
$$

Where $v_{j}$ is the solution of

$$
\begin{align*}
\Delta v_{j} & =0 \quad \text { in } \quad \Omega^{c}, \\
v_{j} & =1 \quad \text { on } \quad \partial \Omega, \tag{2.17}
\end{align*}
$$

With $v_{j} \rightarrow 0$ at infinity. By definition,

$$
\begin{equation*}
\operatorname{cap}(\Omega)=\int \sum\left|\nabla v_{j}\right|^{2} \tag{2.18}
\end{equation*}
$$

Inequality (2.7) still holds but with the constant $k$ replaced by $k=\operatorname{cap}(\Omega)$. Also we note that (2.7) can be written as

$$
\begin{equation*}
\int \sum\left|\nabla f_{j}\right|^{2} \geq S_{3+\epsilon}\left(\sum\left\|f_{j}\right\|_{\frac{2(3+\epsilon)}{1+\epsilon}}^{2}\right)+\sum\left(\int \nabla f_{j} \cdot \nabla u_{j}\right)^{2} /\left[\int\left|\nabla u_{j}\right|^{2}+k\left\|u_{j}\right\|_{\infty}^{2}\right] \tag{2.19}
\end{equation*}
$$

which holds for any $u_{j} \in C_{0}^{\infty}(\Omega)$. By density, (2.19) still holds for every $u_{j}$ in $H_{0}^{1} \cap L^{\infty}$ (the reason is that for every such $u_{j}$ there is a sequence $\left(u_{j}\right)_{j_{0}} \in C_{0}^{\infty}(\Omega)$ with $\left(u_{j}\right)_{j_{0}} \rightarrow u_{j}$ in $H_{0}^{1}$ and $\left.\left\|\left(u_{j}\right)_{j_{0}}\right\|_{\infty} \rightarrow\left\|u_{j}\right\|_{\infty}\right)$.

We now choose $u_{j}$ to be the solution of (2.3) with

$$
\begin{equation*}
\sum g_{j}=\frac{\partial}{\partial x_{i}}\left[\sum\left(\operatorname{sgn} \frac{\partial f_{j}}{\partial x_{i}}\right) 1_{A}\right] \tag{2.20}
\end{equation*}
$$

This function $u_{j}$ is in $L^{\infty}$ as we now verify. We can write

$$
u_{j}=w_{j}+h_{j}
$$

where $w_{j}$ satisfies $\Delta w_{j}=g_{j}$ in all of $\mathbb{R}^{3+\epsilon}$, namely,

$$
\begin{equation*}
w_{j}=C_{3+\epsilon}|x|^{-(1+\epsilon)} * g_{j} \tag{2.21}
\end{equation*}
$$

Clearly $h_{j}$ is harmonic and $h_{j}=-w_{j}$ on $\partial \Omega$ therefore $\left\|\sum h_{j}\right\|_{\infty} \leq\left\|\sum w_{j}\right\|_{\infty, \partial \Omega} \leq\left\|\sum w_{j}\right\|_{\infty}$ and hence $\sum\left\|u_{j}\right\|_{\infty} \leq$ $2 \sum\left\|w_{j}\right\|_{\infty}$. On the other hand, and thus

$$
w_{j}=C_{3+\epsilon} \sum\left(\frac{\partial}{\partial x_{i}}|\mathrm{x}|^{-(1+\epsilon)}\right) *\left[\left(\operatorname{sgn} \frac{\partial f_{j}}{\partial x_{i}}\right) 1_{A}\right],
$$

and thus

$$
\begin{equation*}
\left|w_{j}\right| \leq C_{3+\epsilon}(1+\epsilon)|x|^{-(2+\epsilon)} * 1_{A} . \tag{2.22}
\end{equation*}
$$

Since $|x|^{-(2+\epsilon)} \in L_{w_{j}}^{3+\epsilon / 2+\epsilon}$ we obtain

$$
\begin{equation*}
\left\|\sum u_{j}\right\|_{\infty} \leq 2 \sum\left\|w_{j}\right\|_{\infty} \leq C_{3+\epsilon}^{\prime}|A|^{1 / 3+\epsilon} . \tag{2.23}
\end{equation*}
$$

Next, let us estimate $\int \Sigma\left|\nabla u_{j}\right|^{2}$. Multiplying (2.3) by $u_{j}$ we have

$$
\int \sum\left|\nabla u_{j}\right|^{2}=\int \sum\left(\operatorname{sgn} \partial f_{j} / \partial x_{i}\right) 1_{A}\left(\partial u_{j} / \partial x_{i}\right) \leq\left[\int \sum\left|\nabla u_{j}\right|^{2}\right]^{1 / 2}|A|^{1 / 2}
$$

and thus

$$
\begin{equation*}
\int \sum\left|\nabla u_{j}\right|^{2} \leq|A| \tag{2.24}
\end{equation*}
$$

Finally, since $\sum f_{j}=0$ on $\partial \Omega$,

$$
\begin{equation*}
\int \sum \nabla f_{j} \cdot \nabla u_{j}=-\int \sum f_{j} \Delta u_{j}=\int \sum\left|\partial f_{j} / \partial x_{i}\right| 1_{A} \tag{2.25}
\end{equation*}
$$

Using these estimates (2.19) we find

Since $|A|^{1-(2 / 3+\epsilon)} \leq|\Omega|^{1-(2 / 3+\epsilon)} \leq S_{3+\epsilon}^{-1} \operatorname{cap}(\Omega)$ by Sobolev's inequality applied to the function $\widetilde{v}_{J}=v_{j}$ in $\Omega^{c}$ and $\widetilde{v}_{J}=1$ in $\Omega$. This completes the proof of (1.6) with the constants given in (2.15).

## III. Proofs of (1.7)-(1.9) and Related Matters

Proof of (1.8)(See [13]): Let us define:

$$
\phi_{j}= \begin{cases}f_{j} & \text { in } \quad \Omega  \tag{3.1}\\ w_{j} & \text { in } \Omega^{c}\end{cases}
$$

Where $w_{j}$ is the harmonic function that vanishes at infinity and agrees with $f_{j}$ on $\partial \Omega$. Using $\phi_{j}$ in (1.1) we find:

$$
\begin{equation*}
\int_{\Omega} \sum\left|\nabla f_{j}\right|^{2}+\int_{\Omega^{c}} \sum\left|\nabla w_{j}\right|^{2} \geq S_{3+\epsilon}\left(\sum\left\|f_{j}\right\|_{\frac{2(3+\epsilon)}{2}}^{1+\epsilon}\right) \tag{3.2}
\end{equation*}
$$

On the other hand, we have

$$
\begin{equation*}
\int_{\Omega^{c}} \sum\left|\nabla w_{j}\right|^{2} \sim \sum\left\|f_{j}\right\|_{H^{1 / 2}(\partial \Omega)}^{2} \tag{3.3}
\end{equation*}
$$

This concludes the proof of (1.8).
Proof of (1.7)(See [13]): Now suppose that $f_{j}$ is a constant on $\partial \Omega$. We shall first investigate the case that $\Omega$ is a ball of radius $R$ centered at zero. In this case $w_{j}(x)=f_{j}(\partial \Omega) R^{(3+\epsilon)-2}|x|^{2-(3+\epsilon)}$. Above Inequality (3.2), then yields (1.7) with:

$$
\begin{equation*}
E(\Omega)=\operatorname{cap}(\Omega)=\sigma_{3+\epsilon} R^{1+\epsilon} / 1+\epsilon=\frac{(3+\epsilon)|\Omega|}{1+\epsilon}\left\{\frac{\sigma_{3+\epsilon}}{(3+\epsilon)|\Omega|}\right\}^{2 / 3+\epsilon} \tag{3.4}
\end{equation*}
$$

Furthermore, (1.7) is a strict inequality with this $E(\Omega)$ because the function $\phi_{j}$ is not of the form (1.3). Also, $E(\Omega)$ given by the sharp constant. To see this we apply (1.9) with $f_{j}=\left(f_{j}\right)_{\varepsilon}$, given by (1.3) with $a=1$ and $x+\epsilon=0=$ center of the ball. We have:

$$
\int_{\mathbb{R}^{3+\epsilon}} \sum\left|\nabla\left(f_{j}\right)_{\varepsilon}\right|^{2}=S_{3+\epsilon}\left(\sum\left\|\left(f_{j}\right)_{\varepsilon}\right\|_{\frac{2(3+\epsilon)}{1+\epsilon}, \mathbb{R}^{3+\epsilon}}^{2}\right) .
$$

On the other hand, as $\varepsilon \rightarrow 0$

$$
\begin{align*}
\int_{\mathbb{R}^{3+\epsilon}} \sum\left|\nabla\left(f_{j}\right)_{\varepsilon}\right|^{2}= & \int_{\Omega} \sum\left|\nabla\left(f_{j}\right)_{\varepsilon}\right|^{2}+\int_{\Omega^{c}} \sum\left|\nabla\left(f_{j}\right)_{\varepsilon}\right|^{2} \\
& =\int_{\Omega} \sum\left|\nabla\left(f_{j}\right)_{\varepsilon}\right|^{2}+\operatorname{cap}(\Omega)\left(\sum\left|\left(f_{j}\right)_{\varepsilon}(\partial \Omega)\right|^{2}\right)+o(1) \tag{3.6}
\end{align*}
$$

Here we have to note that as $\varepsilon \rightarrow 0$ for $|x|>R$

$$
\left(f_{j}\right)_{\varepsilon}(x) \rightarrow|x|^{-(1+\epsilon)}
$$

in the appropriate topologies. On the other hand,

$$
\int_{\mathbb{R}^{3+\epsilon}} \sum\left|\left(f_{j}\right)_{\varepsilon}\right|^{\frac{2(3+\epsilon)}{1+\epsilon}}-\int_{\Omega} \sum\left|\left(f_{j}\right)_{\varepsilon}\right|^{\frac{2(3+\epsilon)}{1+\epsilon}}=\int_{\Omega^{c}} \sum\left|\left(f_{j}\right)_{\varepsilon}\right|^{\frac{2(3+\epsilon)}{1+\epsilon}} \rightarrow C
$$

Thus

$$
\begin{equation*}
\sum\left\|\left(f_{j}\right)_{\varepsilon}\right\|_{\frac{2(3+\epsilon)}{1+\epsilon}, \mathbb{R}^{3+\epsilon}}^{2}=\sum\left\|\left(f_{j}\right)_{\varepsilon}\right\|_{\frac{2(3+\epsilon)}{1+\epsilon}, \Omega}^{2}+o(1) \tag{3.7}
\end{equation*}
$$

This proves that $E(\Omega)$ in (1.7) is greater than or equal to $\operatorname{cap}(\Omega)$ when $\Omega$ is a ball, and thus that (3.4) is sharp.
The same calculation with $\left(f_{j}\right)_{\varepsilon}$, as above shows that if $\Omega$ is a ball there is no inequality of the type:

$$
\begin{equation*}
\int_{\Omega} \sum\left|\nabla f_{j}\right|^{2}+\operatorname{cap}(\Omega)\left(\sum\left|f_{j}(\partial \Omega)\right|^{2}\right) \geq S_{3+\epsilon}\left(\sum\left\|f_{j}\right\|_{\frac{2(3+\epsilon)}{1+\epsilon}}^{2}\right)+d \sum\left\|f_{j}\right\|_{1}^{2} \tag{3.8}
\end{equation*}
$$

with $\epsilon \geq 0$, because the additional term $\sum\left\|\left(f_{j}\right)_{\varepsilon}\right\|_{1}=O(1)$ as $\varepsilon \rightarrow 0$.
Now we consider a general domain with $f_{j}(\partial \Omega)=$ constant $=C$. We can assume $C \geq 0$ and note that we can also assume $f_{j} \geq C$ in $\Omega$. (This is so because replacing $f_{j}$ by If $\sum\left|f_{j}-C\right|+C \geq \sum f_{j}$ does not decrease the $L^{\frac{2(3+\epsilon)}{1+\epsilon}}$ norm and leaves $\left\|\sum \nabla f_{j}\right\|_{2}$ invariant.) Consider the function $g_{j}=\sum f_{j}-C \geq 0$ which vanishes on $\partial \Omega$ and hence can be extended to be zero on $\Omega^{c}$. Apply to $g_{j}$ the rearrangement inequality for the $L^{2}$ norm of the
gradient, as was done in Section II. Finally considers $\tilde{f}_{j}=g_{j}^{*}+C$ in the ball $\Omega^{*}$ whose volume is $|\Omega|$. Since $\tilde{f}_{j}\left(\partial \Omega^{*}\right)=C=f_{j}(\partial \Omega)$ we have

$$
\int_{\Omega^{*}} \sum\left|\nabla \tilde{f}_{j}\right|^{2}+E\left(\Omega^{*}\right)\left(\sum\left|f_{j}(\partial \Omega)\right|^{2}\right) \geq S_{n}\left(\sum\left\|\tilde{f}_{j}\right\|_{\frac{2(3+\epsilon)}{2}+\Omega^{*}}^{2+\epsilon}\right),
$$

As we remarked, $\left\|\sum \nabla f_{j}\right\|_{2} \geq\left\|\Sigma \nabla \tilde{f}_{j}\right\|_{2}$. Also since $f_{j} \geq C$, it is easy to check that $\sum\left\|f_{j}\right\|_{\frac{2(3+\epsilon)}{1+\epsilon}}=\sum\left\|\tilde{f}_{j}\right\|_{\frac{2(3+\epsilon)}{1+\epsilon}}$.
The conclusion to be drawn from this exercise is that (1.7) holds for general $\Omega$ with $E(\Omega)$ given by (3.4), namely, $\operatorname{cap}\left(\Omega^{*}\right)$. We also note that (1.7), with this $E(\Omega)$, is strict, since it is strict for a ball.

Question: Is $E(\Omega)$ given by (3.4) the sharp constant in general?
Proof of (1.9)(See [13]): Given $f_{j}$ in $\Omega$ we consider the harmonic function $h_{j}$ in $\Omega$ which equals $f_{j}$ on $\partial \Omega$ We write

$$
\begin{equation*}
f_{j}=h_{j}+u_{j} \tag{3.9}
\end{equation*}
$$

With $\mathrm{u}_{j}=0$ on $\partial \Omega$ and thus

$$
\begin{equation*}
\int \sum\left|\nabla u_{j}\right|^{2} \geq S_{3+\epsilon}\left(\sum\left\|u_{j}\right\|_{\left.\frac{2(3+\epsilon)}{1+\epsilon}\right)}^{2}\right) \tag{3.10}
\end{equation*}
$$

On the one hand

$$
\begin{equation*}
\int \sum\left|\nabla u_{j}\right|^{2}=\int \sum\left|\nabla\left(f_{j}-h_{j}\right)\right|^{2}=\int \sum\left|\nabla f_{j}\right|^{2}-\int \sum\left|\nabla h_{j}\right|^{2} \tag{3.11}
\end{equation*}
$$

(note that $\int_{\Omega} \Sigma\left|\nabla h_{j}\right|^{2}=\int_{\partial \Omega} \sum h_{j}\left(\partial h_{j} / \partial 3+\epsilon\right)=\int_{\partial \Omega} \sum f_{j}\left(\partial h_{j} / \partial 3+\epsilon\right)=\int_{\Omega} \Sigma\left(\nabla f_{j} \nabla h_{j}\right)$ ). On the other hand, by the triangle inequality,

$$
\begin{equation*}
\sum\left\|u_{j}\right\|_{\frac{2(3+\epsilon)}{1+\epsilon}} \geq \sum\left\|f_{j}\right\|_{\frac{2(3+\epsilon)}{1+\epsilon}}-\sum\left\|h_{j}\right\|_{\frac{2(3+\epsilon)}{1+\epsilon}} . \tag{3.12}
\end{equation*}
$$

Inserting (3.11) and (3.12) in (3.10) we obtain

$$
\begin{equation*}
\sum\left\|\nabla f_{j}\right\|_{2}+\sum\left\|h_{j}\right\|_{\frac{2(3+\epsilon)}{1+\epsilon}} \geq S_{3+\epsilon}^{1 / 2}\left(\sum\left\|f_{j}\right\|_{\left.\frac{2(3+\epsilon)}{1+\epsilon}\right)}\right) \tag{3.13}
\end{equation*}
$$

Next we claim that

$$
\begin{equation*}
\sum\left\|h_{j}\right\|_{\frac{2(3+\epsilon)}{1+\epsilon}} \leq G(\Omega)\left(\sum\left\|f_{j}\right\|_{\frac{3+\epsilon}{2+\epsilon}, \partial \Omega}\right) \tag{3.14}
\end{equation*}
$$

with $\epsilon^{2}+4 \epsilon+5=0$, which will complete the proof of (1.9). The proof is a standard duality argument. Indeed, let $\psi_{j}$ be the solution of

$$
\begin{array}{rll}
\Delta \psi_{j}=Y & \text { in } \quad \Omega \\
\psi_{j}=0 & \text { on } & \partial \Omega \tag{3.15}
\end{array}
$$

where $Y$ is some arbitrary function in $L^{t}$. We have, by multiplying by $h_{j}$ and integrating by parts,

$$
\begin{equation*}
\int_{\Omega} \sum h_{j} Y=\int_{\partial \Omega} \sum f_{j} \frac{\partial \psi_{j}}{\partial(3+\epsilon)} \tag{3.16}
\end{equation*}
$$

However, the $L^{\frac{3+\epsilon}{1+\epsilon}}$ regularity theory shows that $\psi_{j} \in W^{2, t}$ with $\left\|\sum \psi_{j}\right\|_{W^{2, t}(\Omega)} \leq C\|Y\|_{t}$. In particular, $\left\|\Sigma \nabla \psi_{j}\right\|_{W^{1, t}(\Omega)} \leq C\|Y\|_{t}$ and, by trace inequalities,

$$
\begin{equation*}
\left\|\sum \frac{\partial \psi_{j}}{\partial 3+\epsilon}\right\|_{\frac{t(2+\epsilon)}{(3+\epsilon)-t^{\prime}} \partial \Omega} \leq C\|Y\|_{t} \tag{3.17}
\end{equation*}
$$

Therefore, by (3.16) and Hölder's inequality,

$$
\begin{equation*}
\left|\int \sum h_{j} Y\right| \leq C \sum\left\|f_{j}\right\|_{\frac{3+\epsilon}{2+\epsilon}, \partial \Omega}\|Y\|_{t} \tag{3.19}
\end{equation*}
$$

Since (3.19) holds for all $Y$ we conclude that

$$
\left\|\sum h_{j}\right\|_{2^{*}} \leq C \sum\left\|f_{j}\right\|_{\frac{3+\epsilon}{2+\epsilon} \partial \Omega^{\prime}}
$$

when $\epsilon^{2}+4 \epsilon+5=0$.
Finally, we claim that there is no inequality of the type (1.9) with $\epsilon^{2}+4 \epsilon+5=0$. Indeed, suppose (1.9) holds with some such $\frac{3+\epsilon}{2+\epsilon}$. We choose $f_{j}=\left(f_{j}\right)_{\varepsilon}$, as in (1.3) with $a=1$ and $(x+\epsilon) \in \partial \Omega$. It is obvious that as $\varepsilon \rightarrow 0$

$$
\begin{gathered}
\sum \int_{\Omega}\left|\nabla\left(f_{j}\right)_{\varepsilon}\right|^{2} / \int_{\mathbb{R}^{3+\epsilon}}\left|\nabla\left(f_{j}\right)_{\varepsilon}\right|^{2}=1 / 2+o(1) \\
\sum \int_{\Omega}\left|\left(f_{j}\right)_{\varepsilon}\right|^{\frac{2(3+\epsilon)}{1+\epsilon}} / \int_{\mathbb{R}^{3+\epsilon}}\left|\left(f_{j}\right)_{\varepsilon}\right|^{\frac{2(3+\epsilon)}{1+\epsilon}}=1 / 2+o(1)
\end{gathered}
$$

while

$$
\begin{aligned}
& \int_{\mathbb{R}^{3+\epsilon}} \sum\left|\nabla\left(f_{j}\right)_{\varepsilon}\right|^{2}=S_{3+\epsilon}\left(\sum\left\|\left(f_{j}\right)_{\varepsilon}\right\|_{\frac{2(3+\epsilon)}{1+\epsilon}, \mathbb{R}^{3+\epsilon}}^{2}\right) \\
& \text { and } \sum\left\|\left(f_{j}\right)_{\varepsilon}\right\|_{\frac{3+\epsilon}{2+\epsilon}, \partial \Omega} /\left\|\left(f_{j}\right)_{\varepsilon}\right\|_{\frac{2(3+\epsilon)}{1+\epsilon}}=o(1)
\end{aligned}
$$

This contradicts (1.9).
Remark. The last exercise with $\left(f_{j}\right)_{\varepsilon}$ given above shows that it is not possible to apply rearrangement techniques when $f_{j}$ is not constant on $\partial \Omega$, even if $\Omega$ is a ball. It also shows that there is no inequality for all $f_{j} \in H^{1}$ of the type

$$
\left\|\sum \nabla f_{j}\right\|_{2}^{2}+C \sum\left\|f_{j}\right\|_{\frac{3+\epsilon}{2+\epsilon^{N}}}^{2} \geq S_{3+\epsilon}\left(\sum\|f\|_{\frac{2(3+\epsilon)}{1+\epsilon}}^{2}\right)
$$

with $\epsilon>-3$.
Proof of (1.11)(See [13]): Let $\Omega$ be a ball of radius $R$ centered at zero. For simplicity, assume $R=1$. Define

$$
g_{j}(x)= \begin{cases}f_{j}(x), & |x| \leq 1  \tag{3.20}\\ |x|^{-(1+\epsilon)} f_{j}\left(x|x|^{-2}\right) & |x| \geq 1\end{cases}
$$

and apply the usual Sobolev inequality (1.1) to $g_{j}$. We note (by a change of variables) that

$$
\begin{align*}
& \int_{\Omega} \sum g_{j}^{\frac{2(3+\epsilon)}{1+\epsilon}}=\int_{\Omega^{c}} \sum g_{j}^{\frac{2(3+\epsilon)}{1+\epsilon}} \\
& \int_{\Omega} \sum\left|\nabla g_{j}\right|^{2}=\int_{\Omega^{c}} \sum\left|\nabla g_{j}\right|^{2}-(1+\epsilon)\left\|f_{j}\right\|_{2, \partial \Omega}^{2} \tag{3.21}
\end{align*}
$$

Inserting (3.21) into (1.1) yields (1.11) with $I(\Omega)=(1+\epsilon) / 2$.

## Remark on the Hardy-Littlewood-Sobolev Inequality

Consider the inequality (in $\mathbb{R}^{3}$ )

$$
\begin{equation*}
\sum I\left(f_{j}\right) \leq P\left(\sum\left\|f_{j}\right\|_{6,5}^{2}\right) \tag{3.22}
\end{equation*}
$$

with

$$
\begin{equation*}
\sum_{i \neq 1} I\left(f_{j}\right)=\iint \sum \mathrm{f}_{j}(x) f_{j}(x+\epsilon)|\epsilon|^{-1} d x d(x+\epsilon) \geq 0 \tag{3.23}
\end{equation*}
$$

The sharp constant $P$ is known to be [7]

$$
\begin{equation*}
P=4^{5 / 3} /\left[3 \pi^{1 / 3}\right] . \tag{3.24}
\end{equation*}
$$

Let $\Omega$ be a ball of radius one centered at zero and assume that $\sum f_{j}=0$ outside $\Omega$. In this case, (3.22) is strict because the only functions that give equality in (3.22) are of the form [7]

$$
\begin{equation*}
\sum\left(f_{j}\right)_{\varepsilon}(x)=a\left[\varepsilon^{2}+|\epsilon|^{2}\right]^{-5 / 2} \tag{3.25}
\end{equation*}
$$

For $\sum f_{j}=0$ outside $\Omega$, we ask whether (3.22) can be improved to

$$
\begin{equation*}
C\left(\sum\left\|f_{j}\right\|_{1}^{2}\right)+\sum I\left(f_{j}\right) \leq P\left(\sum\left\|f_{j}\right\|_{6 / 5}^{2}\right) \tag{3.26}
\end{equation*}
$$

Our conclusion is that (3.26) fails for any $C>0$.
Take $f_{j}=\left(\widetilde{f}_{J}\right)_{\varepsilon}=\left(f_{j}\right)_{\varepsilon} 1_{\Omega}$ with $\left(f_{j}\right)_{\varepsilon}$ given by (3.25) and with $x+\epsilon=0$ and with $a=a_{\varepsilon}$ chosen so that $\sum\left\|\left(f_{j}\right)_{\varepsilon}\right\|_{6 / 5, \mathbb{R}^{3}}=1$. The function $\left(f_{j}\right)_{\varepsilon}$ satisfies the following (Euler) equation on $\mathbb{R}^{3}$,

$$
\begin{equation*}
\sum \frac{1}{|x|} *\left(f_{j}\right)_{\varepsilon}=P\left(\sum\left(f_{j}\right)_{\varepsilon}^{1 / 5}\right) \tag{3.27}
\end{equation*}
$$

However, for $|x|<1$

$$
\begin{equation*}
\sum\left(\frac{1}{|x|} *\left(\widetilde{f}_{J}\right)_{\varepsilon}\right)(x)+K_{\varepsilon}=\sum\left(\frac{1}{|x|} *\left(f_{j}\right)_{\varepsilon}\right)(x), \tag{3.28}
\end{equation*}
$$

where $K_{\varepsilon}$ is a constant bounded above by $D_{\varepsilon}=\int_{|x|>1} \sum\left(f_{j}\right)_{\varepsilon}$. Multiply (3.27) by $\left(\widetilde{f}_{J}\right)_{\varepsilon}$ and integrate over $\Omega$. Then

$$
\begin{align*}
& \sum I\left(\widetilde{f}_{J}\right)_{\varepsilon}+T_{\varepsilon}\left(\sum\left\|\left(\widetilde{f}_{J}\right)_{\varepsilon}\right\|_{1}^{2}\right) \geq \sum I\left(\widetilde{f}_{J}\right)_{\varepsilon}+K_{\varepsilon} \int \sum\left(\widetilde{f}_{J}\right)_{\varepsilon} \\
&=P\left(\sum\left\|\left(\widetilde{f}_{J}\right)_{\varepsilon}\right\|_{6 / 5}^{6 / 5}\right) \geq \mathrm{P}\left(\sum\left\|\left(\widetilde{f}_{J}\right)_{\varepsilon}\right\|_{6 / 5}^{2}\right) \tag{3.29}
\end{align*}
$$

where $\mathrm{T}_{\varepsilon}=\mathrm{D}_{\varepsilon} / \int \sum\left(\tilde{\mathrm{f}}_{\mathrm{J}}\right)_{\varepsilon}$. From (3.29), we see that (3.26) fails if $\mathrm{C}>\mathrm{T}_{\varepsilon}$ for any $\varepsilon>0$. However, it is obvious that $\mathrm{T}_{\varepsilon} \rightarrow 0$ as $\varepsilon \rightarrow 0$.

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