

Oscillatory Magnetopolar Free Convection Flow Through a Vertical Porous Plate Embedded in a Porous Medium in Slip Flow Regime

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Abstract: *In this paper we study as Oscillatory two dimensional magnetopolar free convection flow through a porous medium with combined heat and mass transfer and thermal radiation in slip flow regime. The permeability and suction velocity are assumed to be time dependent. Using perturbation technique expressions for velocity (u), angular velocity (ω), temperature (θ), concentration (C), skin friction (C_f) and Nusselt number (Nu) are obtained and a comparative study is made to analyze the effects of different parameters. We notice that as we increase permeability parameter (K) skin friction falls in the beginning but rises as we move away from the plate. Moreover, for both the basic fluids air ($Pr = 0.71$, $Sc = 0.22$) and water ($Pr = 7$, $Sc = 0.61$), velocity increases on decreasing the slip at the boundary.*

Key Words: *Heat flux, Mass flux, Polar fluid, Radiation, Unsteady.*

I. Introduction

The natural or free convection process is present in various physical phenomenon such as fire engineering, nuclear energy, fiber and granular insulation, geothermal system etc. Simultaneously, heat and mass transfer from different geometrics embedded in porous media have many engineering and geophysical applications such as geothermal reservoirs, drying of porous solids, thermal insulation, enhanced oil recovery etc. Magnetohydrodynamics has attracted the attention of a large number of scholars due to its diverse applications. In engineering it finds its applications in MHD pumps, MHD bearings etc. The phenomenon of mass transfer is also very common in theory of stellar structure and observable effects are detectable, at least on the solar surface. The study of effects of magnetic field on free convection flow is important in liquid metals, electrolyte and ionized gases.

Das and Das [1] have studied MHD free convection flow near a moving plate in presence of thermal radiation where as MHD free convection flow of a viscous fluid through a porous medium bounded by an oscillating plate was studied by Singh and Gupta [2]. Moreover, Ahmed [3] observed the effects of unsteady free convection MHD flow through a porous medium while Sahoo et al. [4] studied MHD unsteady free convection flow with constant suction and heat sink. Cookey et al. [5] studied the influence of viscous dissipation and radiation on unsteady MHD free convection flow in a porous medium with time dependent suction on the other hand Aboeldahab and Azzam [6] made studies on unsteady three dimensional combined heat and mass transfer with time dependent chemical reaction.

Convection heat transfer through a porous medium has been a subject of great interest for the last few decades as these are quite prevalent in nature.

Such flows have many engineering applications in petroleum technology to study the movement of natural gas, oil and water through oil reservoirs, in chemical engineering for filtration and purification process etc. In view of these applications, many scholars have made investigations where porous medium is either bounded by horizontal or vertical surface. Magyarai et al. [7] found an analytic solution for unsteady free convection in porous media where as Geindreau and Auriault [8] studied MHD flows in porous media. Ahmed and Ahmed [9, 10] studied oscillatory two dimensional as well as three dimensional flows through a porous medium with viscous dissipative heat.

The Navier-stokes model of classical hydrodynamics has drastic limitations. It cannot describe fluids with microstructure, fluids that are interesting in themselves and important in application. In general, individual particles of such complex fluids, eg. Polymeric suspensions, animal blood, liquid crystals etc., may be of different shapes, may shrink and expand, or change their shape and moreover they may rotate independently of the rotation and movement of the fluid. Such fluids are called Polar fluids which belong to the class of fluids with non-symmetric stress tensor. These fluids are more general than that considered by classical hydrodynamics that we call ordinary fluids {Lukaszewicz [11]}. Kim [12] studied unsteady MHD convection flow of a polar fluid in a porous medium Patil and Kulkarni [13] studied effects of chemical reaction on free convection flow of a polar fluid. Jain and Gupta [14, 15] studied effects of rotational parameter on unsteady magnetopolar free convection

with different boundary conditions while Cheng [16] studied natural convection heat and mass transfer from a sphere in micropolar fluids.

It is a well known fact that in case of many polymeric liquids when the weight of the molecule is high, it shows slip at the boundary. In many problems like that of thin films, rarefied fluids, fluid containing concentrated suspension, the no slip boundary condition fails to work. Jain et al. [17] and Chaudhary and Jain [18] have studied effects of different parameters on combined heat and mass transfer in magneto polar fluid with slip flow regime. Fang et al. [19] studied slip MHD viscous flow over a stretching sheet while Sahin [20] studied the influence of chemical reaction on transient MHD free convective flow in slip flow regime.

In this paper we study an oscillatory magnetopolar free convection flow through a porous medium with time dependent permeability and suction velocity. The flow is also effected by rotational and couple stress parameters with combined heat and mass transfer and thermal radiation with slip flow regime. Using perturbation technique solutions have been obtained for velocity (u), angular velocity (ω) temperature (T), concentration (C), skin friction (C_f) and Nusselt number (Nu) and are shown graphically for both the basic fluids air (Pr = 0.71, Sc = 0.22) and water (Pr = 7, Sc = 0.61).

II. Formulation of the problem

An unsteady MHD free convective, two dimensional heat and mass transfer flow with radiation of a polar fluid through a porous medium in a slip flow regime with transverse magnetic field of strength B_0 is considered. The plate is subjected to a constant heat and mass flux. The suction velocity and permeability of the porous medium are assumed to time dependent and are of the form as shown in the figure:

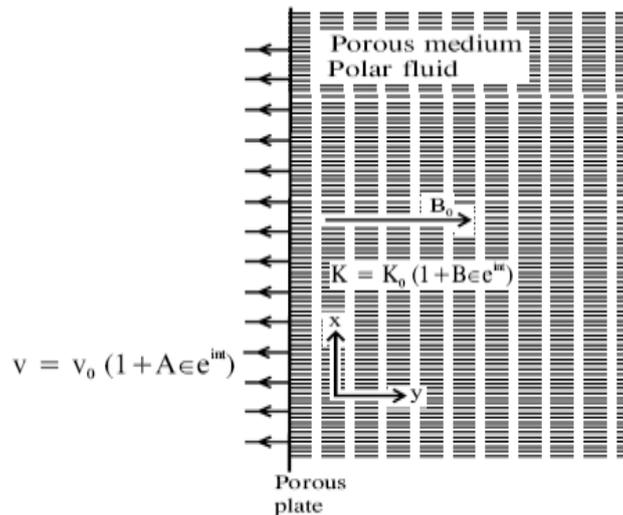


Figure : Schematic diagram.

where K_0 is the mean permeability of the medium, v_0 is the mean suction velocity, n is the frequency of fluctuation, t the time and $\omega \ll 1$, A and B are constants.

The axis of x is taken along the plate in vertically upward direction and y -axis is taken normal to the plate. Under these conditions and using the Boussineque's approximation, governing equations of the flow are given by:

$$\frac{\partial u}{\partial t} - v_0(1 + A e^{int}) \frac{\partial u}{\partial y} = g\beta(T - T_\infty) + g\beta^*(C - C_\infty) + (v + v_r) \frac{\partial^2 u}{\partial y^2} + 2v_r \frac{\partial \omega}{\partial y} - \frac{v}{K_0(1 + B e^{int})} u - \frac{\sigma B_0^2}{\rho} u \tag{1}$$

$$\frac{\partial \omega}{\partial t} - v_0(1 + A e^{int}) \frac{\partial \omega}{\partial y} = \frac{\gamma}{I} \frac{\partial^2 \omega}{\partial y^2} \tag{2}$$

$$\frac{\partial T}{\partial t} - v_0(1 + A e^{int}) \frac{\partial T}{\partial y} = \frac{\kappa}{\rho C_p} \frac{\partial^2 T}{\partial y^2} - \frac{1}{\rho C_p} \frac{\partial q_r}{\partial y} \tag{3}$$

$$\frac{\partial C}{\partial t} - v_0(1 + A \epsilon e^{int}) \frac{\partial C}{\partial y} = D \frac{\partial^2 C}{\partial y^2} \tag{4}$$

where u, ω, T and C are velocity, angular velocity, temperature and species concentration respectively of the fluid particles, g is acceleration due to gravity, β is coefficient of volume expansion, β_c is coefficient of species concentration expansion, $\rho, \nu, \nu_r, C_p, B, D, K$ are density, kinematic viscosity, rotational kinematic viscosity, thermal conductivity, specific heat at constant pressure, electrical conductivity, magnetic field, mass diffusivity and permeability of porous medium respectively. I a scalar constant equal to moment of Inertia of unit mass and

$$\gamma = C_a + C_d$$

where C_a and C_d are coefficient of couple stress viscosities.

The local radiant heat for the case of an optically thin grey gas in expressed by England and Emery [21] as

$$\frac{\partial q_r}{\partial y} = -4a^* \sigma^* (T_\infty^4 - T^4),$$

where $T^4 \cong 4T_\infty^3 T - 3T_\infty^4$

σ^* is Stefan-Boltzmann constant and a^* is absorption coefficient.

The boundary conditions are:

$$\left. \begin{aligned} y = 0 : u = L_1 \frac{\partial u}{\partial y}, \frac{\partial \omega}{\partial y} = -\frac{\partial^2 u}{\partial y^2}, \frac{\partial T}{\partial y} = -\frac{q}{\kappa}, \frac{\partial C}{\partial y} = -\frac{m}{D} \\ y \rightarrow \infty : u \rightarrow 0, \quad \omega \rightarrow 0, \quad T \rightarrow T_\infty, \quad C \rightarrow C_\infty \end{aligned} \right\} \tag{5}$$

where $L_1 = (2 - m_1/m_1)L$, m_1 being the Maxwell's reflection coefficient and L the free path.

On introducing the following non-dimensional quantities:

$$y^* = \frac{y v_0}{\nu}, \quad t^* = \frac{t v_0^2}{4\nu}, \quad n^* = \frac{4\nu n}{v_0^2}$$

$$u^* = \frac{u}{v_0}, \quad \omega^* = \frac{\nu \omega}{v_0^2},$$

$$\theta = \frac{(T - T_\infty) \kappa v_0}{\nu q}, \quad C^* = \frac{(C - C_\infty) D v_0}{\nu m},$$

$$K^* = \frac{K_0 v_0^2}{\nu^2} \text{ (permeability parameter),}$$

equations (1) to (4) in non-dimensional form, after dropping the asterisks over them, reduces to

$$\frac{1}{4} \frac{\partial u}{\partial t} - (1 + A \epsilon e^{int}) \frac{\partial u}{\partial y} = Gr\theta + GcC + (1 + \alpha_1) \frac{\partial^2 u}{\partial y^2} + 2\alpha_1 \frac{\partial \omega}{\partial y} - \left\{ M + \frac{1}{K(1 + B \epsilon e^{int})} \right\} u \tag{6}$$

$$\frac{1}{4} \frac{\partial \omega}{\partial t} - (1 + A \epsilon e^{int}) \frac{\partial \omega}{\partial y} = \frac{1}{\beta_1} \frac{\partial^2 \omega}{\partial y^2} \tag{7}$$

$$\frac{1}{4} \frac{\partial \theta}{\partial t} - (1 + A \epsilon e^{int}) \frac{\partial \theta}{\partial y} = \frac{1}{Pr} \frac{\partial^2 \theta}{\partial y^2} - \frac{R}{Pr} \theta \tag{8}$$

$$\frac{1}{4} \frac{\partial C}{\partial t} - (1 + A \epsilon e^{int}) \frac{\partial C}{\partial y} = \frac{1}{Sc} \frac{\partial^2 C}{\partial y^2} \tag{9}$$

where

$$h_1 = \frac{v_0 L_1}{\nu} \text{ (velocity slip parameter),}$$

$$Gr = \frac{g\beta v_0^2 q}{\kappa v_0^4} \text{ (thermal Grashof number),}$$

$$Gc = \frac{g\beta^* v_0^2 m}{D v_0^4} \text{ (mass Grashof number),}$$

$$M = \frac{\sigma B_0^2 \nu}{\rho v_0^2} \text{ (magnetic parameter),}$$

$$\alpha_1 = \frac{\nu_r}{\nu} \text{ (rotational parameter),}$$

$$\beta_1 = \frac{I\nu}{\gamma} \text{ (couple stress parameter),}$$

$$Pr = \frac{\mu C_p}{\kappa} \text{ (Prandtl number),}$$

$$Sc = \frac{\nu}{D} \text{ (Schmidt number),}$$

$$R = \frac{16a^* \sigma^* v_0^2 T_\infty^3}{v_0^2 \kappa} \text{ (radiation parameter).}$$

With corresponding boundary conditions as:

$$\left. \begin{aligned} y = 0 : u = h_1 \frac{\partial u}{\partial y}, \quad \frac{\partial \omega}{\partial y} = -\frac{\partial^2 u}{\partial y^2}, \quad \frac{\partial \theta}{\partial y} = -1, \quad \frac{\partial C}{\partial y} = -1 \\ y \rightarrow \infty : u \rightarrow 0, \quad \omega \rightarrow 0, \quad \theta \rightarrow 0, \quad C \rightarrow 0 \end{aligned} \right\} \quad (10)$$

III. Solution of the problem

Since $\epsilon \ll 1$ is very small and to reduce the system of partial differential equations (6) to (9) into ordinary differential equations, we assume:

$$f(y, t) = f_0(y) + \epsilon \text{int} f_1(y) \quad (11)$$

Where f stands for u, ω, θ and C .

Substituting expression (11) in (6) to (9) and comparing the coefficients of identical powers of ϵ we obtain

$$\begin{aligned} u_0'' + \frac{1}{(1+\alpha_1)} u_0' - \frac{\left(M + \frac{1}{K}\right)}{(1+\alpha_1)} u_0 = -\frac{Gr}{(1+\alpha_1)} \theta_0 - \frac{Gc}{(1+\alpha_1)} C_0 \\ - \frac{2\alpha_1}{(1+\alpha_1)} \omega_0' \end{aligned} \quad (12)$$

$$\begin{aligned} u_1'' + \frac{1}{(1+\alpha_1)} u_1' - \frac{\left(M + \frac{1}{K} + \frac{in}{4}\right)}{(1+\alpha_1)} u_1 = -\frac{Gr}{(1+\alpha_1)} \theta_1 - \frac{Gc}{(1+\alpha_1)} C_1 \\ - \frac{A}{(1+\alpha_1)} u_0' - \frac{2\alpha_1}{(1+\alpha_1)} \omega_1' - \frac{B}{K(1+\alpha_1)} \end{aligned} \quad (13)$$

$$\omega_0'' + \beta_1 \omega_0' = 0 \quad (14)$$

$$\omega_0'' + \beta_1 \omega_1' - \frac{in\beta_1}{4} \omega_1 = -A\beta_1 \omega_0' \quad (15)$$

$$\theta_0'' + Pr \theta_0' - R \theta_0 = 0 \quad (16)$$

$$\theta_1'' + \text{Pr} \theta_1' - \left(R + \frac{\text{Pr} \text{in}}{4} \right) \theta_1 = -A \text{Pr} \theta_0' \quad (17)$$

$$C_0'' + \text{Sc} C_0' = 0 \quad (18)$$

$$C_1'' + \text{Sc} C_1' - \frac{\text{Sc} \text{in}}{4} C_1 = -A \text{Sc} C_0' \quad (19)$$

with boundary conditions:

$$\left. \begin{aligned} y = 0 : u_0 &= h_1 u_0', \omega_0' = -u_0'', \theta_0' = -1, C_0' = -1 \\ u_1 &= h_1 u_1', \omega_1' = -u_1'', \theta_1' = 0, C_1' = 0 \\ y \rightarrow \infty : u_0 &\rightarrow 0, \omega_0 \rightarrow 0, \theta_0 \rightarrow 0, C_0 \rightarrow 0 \\ u_1 &\rightarrow 0, \omega_1 \rightarrow 0, \theta_1 \rightarrow 0, C_1 \rightarrow 0 \end{aligned} \right\} \quad (20)$$

Solving (12) to (19) under boundary conditions (20) we get:

$$u_0 = m_7 e^{x_9 y} + J_3 e^{x_3 y} + J_u \bar{e}^{\text{Sc} y} + J_5 \bar{e}^{\beta_1 y} \quad (21)$$

$$u_1 = m_8 e^{x_{11} y} + J_7 e^{x_5 y} + J_8 e^{x_7 y} + J_9 e^{x_{11} y} + J_{10} e^{x_3 y} + J_{11} \bar{e}^{\text{Sc} y} + J_{12} \bar{e}^{\beta_1 y} + J_{13} e^{x_9 y} \quad (22)$$

$$\omega_0 = m_1 \bar{e}^{\beta_1 y} \quad (23)$$

$$\omega_1 = m_2 e^{x_1 y} + J_6 \bar{e}^{\beta_1 y} \quad (24)$$

$$\theta_0 = m_3 e^{x_3 y} \quad (25)$$

$$\theta_1 = m_4 e^{x_5 y} + J_1 e^{x_3 y} \quad (26)$$

$$C_0 = m_5 \bar{e}^{\text{Sc} y} \quad (27)$$

$$C_1 = m_6 e^{x_7 y} + J_2 \bar{e}^{\text{Sc} y} \quad (28)$$

Now substituting equations (21) to (28) in (11) and separating the real and imaginary parts, we get the expressions of real part for velocity, angular velocity, temperature and concentration as:

$$u = u_0(y) + \epsilon (M r_1 \cos nt - M i_1 \sin nt) \quad (29)$$

$$\omega = \omega_0(y) + \epsilon (M r_2 \cos nt - M i_2 \sin nt) \quad (30)$$

$$\theta = \theta_0(y) + \epsilon (M r_3 \cos nt - M i_3 \sin nt) \quad (31)$$

$$C = C_0(y) + \epsilon (M r_4 \cos nt - M i_4 \sin nt) \quad (32)$$

where

$$M r_1 = d_{66} + d_{16} + d_{18} + d_{20} + d_{22} + d_{24} + d_{26}$$

$$M i_1 = d_{67} + d_{17} + d_{19} + d_{21} + d_{23} + d_{25} + d_{27} + d_{28}$$

$$M r_2 = e^{d_8 y/2} (d_6 C_1 - d_7 C_2)$$

$$M i_2 = e^{d_8 y/2} (d_6 C_2 + d_7 C_1) + \frac{4A u_0''(0)}{n} e^{-\beta_1 y}$$

$$M r_3 = d_{34} + d_{36}, \quad M i_3 = d_{35} + d_{37}$$

$$M r_4 = d_{42}, \quad M i_4 = d_{43} + d_{44}$$

IV. Skin Friction

When velocity component u is obtained, now we can now obtain the important parameter viz. skin friction in real part as:

$$C_f = \frac{\tau_w}{\rho v_0^2} = (1 + \alpha_1) \left\{ (m_7 x_9 + J_3 x_3 - J_4 Sc - J_5 \beta_1) + \in (t_{16} \cos nt - t_{17} \sin nt) \right\} \quad (33)$$

V. Nusselt Number

Another important physical parameter of interest viz. Nusselt number in dimensionless form is:

$$Nu = \frac{1}{\theta(0)}$$

where

$$\theta(0) = m_3 + \in \left\{ (t_{18} \cos nt - t_{19} \sin nt) + i(t_{18} \sin nt + t_{19} \cos nt) \right\}.$$

Taking the real part only, we get:

$$Nu = \frac{t_{20}}{t_{20}^2 + t_{21}^2}. \quad (34)$$

Where

$$\begin{aligned} x_1 &= \frac{-\beta_1 - \sqrt{\beta_1^2 + in\beta_1}}{2}, & x_2 &= \frac{-\beta_1 + \sqrt{\beta_1^2 + in\beta_1}}{2} \\ x_3 &= \frac{-Pr - \sqrt{Pr^2 + 4R}}{2}, & x_4 &= \frac{-Pr + \sqrt{Pr^2 + 4R}}{2} \\ x_5 &= \frac{-Pr - \sqrt{Pr^2 + 4\left(R + \frac{Prin}{4}\right)}}{2}, & x_6 &= \frac{-Pr + \sqrt{Pr^2 + 4\left(R + \frac{Prin}{4}\right)}}{2} \\ x_7 &= \frac{-Sc - \sqrt{Sc^2 + Scin}}{2}, & x_8 &= \frac{-Sc + \sqrt{Sc^2 + Scin}}{2} \\ x_9 &= \frac{-1 - \sqrt{1 + 4\left(M + \frac{1}{K}\right)(1 + \alpha_1)}}{2(1 + \alpha_1)}, & x_{10} &= \frac{-1 + \sqrt{1 + 4\left(M + \frac{1}{K}\right)(1 + \alpha_1)}}{2(1 + \alpha_1)} \\ x_{11} &= \frac{-1 - \sqrt{1 + 4\left(M + \frac{1}{K} + \frac{in}{4}\right)(1 + \alpha_1)}}{2(1 + \alpha_1)}, & x_{12} &= \frac{-1 + \sqrt{1 + 4\left(M + \frac{1}{K} + \frac{in}{4}\right)(1 + \alpha_1)}}{2(1 + \alpha_1)} \\ m_1 &= \frac{u_0''(0)}{\beta_1}, & m_2 &= \frac{1}{x_1} [-u_1''(0) + J_6 \beta_1] \\ m_3 &= -\frac{1}{x_3}, & m_4 &= \frac{-Ax_3}{x_5(x_3 - x_6)(x_3 - x_5)} \\ m_5 &= \frac{1}{Sc}, & m_6 &= \frac{A Sc}{x_7(Sc + x_7)(Sc + x_8)} \\ m_7 &= \frac{J_3(h_1 x_{3-1}) - J_4(1 + h_1 Sc) - J_5(1 + h_1 \beta_1)}{(1 - h_1 x_9)} \end{aligned}$$

$$\begin{aligned}
 J_3 &= \frac{-Gr m_3}{(1 + \alpha_1)(x_3 - x_9)(x_3 - x_{10})}, & J_4 &= \frac{-Gc m_5}{(1 + \alpha_1)(Sc + x_{10})(Sc + x_9)} \\
 J_5 &= \frac{2\alpha_1 m_1 \beta_1}{(1 + \alpha_1)(\beta_1 + x_{10})(\beta_1 + x_9)}, & J_1 &= \frac{A Pr}{(x_3 - x_5)(x_3 - x_6)} \\
 J_2 &= \frac{A Sc}{(Sc + x_7)(Sc + x_8)}, & J_6 &= \frac{A m_1 \beta_1^2}{(\beta_1 + x_2)(\beta_1 + x_1)} \\
 J_7 &= \frac{-Gr m_4}{(1 + \alpha_1)(x_5 - x_{12})(x_5 - x_{11})}, & J_8 &= \frac{-Gc m_6}{(1 + \alpha_1)(x_7 - x_{12})(x_7 - x_{11})} \\
 J_9 &= \frac{-2\alpha_1 m_2 x_1}{(1 + \alpha_1)(x_1 - x_{11})(x_1 - x_{12})}, & J_{10} &= \frac{b_1}{(1 + \alpha_1)(x_3 - x_{11})(x_3 - x_{12})} \\
 J_{11} &= \frac{b_2}{(1 + \alpha_1)(Sc + x_{11})(Sc + x_{12})}, & J_{12} &= \frac{b_3}{(1 + \alpha_1)(\beta_1 + x_{11})(\beta_1 + x_{12})} \\
 J_{13} &= \frac{b_4 m_7}{(1 + \alpha_1)(x_9 - x_{11})(x_9 - x_{12})}, & b_1 &= \left(-Gr J_1 - A J_3 x_3 - \frac{J_3}{K} \right) \\
 b_2 &= \left(-Gc J_2 + A J_4 Sc - \frac{J_4}{K} \right), & b_3 &= 2\alpha_1 \beta_1 J_6 + A J_5 \beta_1 - \frac{J_5}{K} \\
 b_4 &= -A x_9 - \frac{B}{K} \\
 A_1 &= \left(\frac{Pr n}{4} \right) \left(\frac{(Pr(1 + \alpha_1) - 1)}{2} \sqrt{\frac{\sqrt{(Pr^2 + 4R)^2 + Pr^2 n^2} - (Pr^2 + 4R)}{2}} + \frac{(Pr(1 + \alpha_1) - 1)n}{4} \right) \\
 B_1 &= \left(-\frac{Pr n}{4} \right) \left((1 + \alpha_1) \frac{Pr^2}{2} + (1 + \alpha_1) R - \frac{Pr}{2} - M - \frac{1}{K} + \frac{(Pr(1 + \alpha_1) - 1)}{2} \right. \\
 &\quad \left. \sqrt{\frac{\sqrt{(Pr^2 + 4R)^2 + Pr^2 n^2} + (Pr^2 + 4R)}{2}} \right) \\
 A_2 &= \left(\frac{n}{4} \right) \left(\frac{(Sc(1 + \alpha_1) - 1)}{2} \sqrt{\frac{Sc \sqrt{Sc^2 + n^2} - Sc^2}{2}} + \frac{(Sc(1 + \alpha_1) - 1)n}{4} \right) \\
 B_2 &= \left(-\frac{n}{4} \right) \left((1 + \alpha_1) \frac{Sc^2}{2} - \frac{Sc}{2} - M - \frac{1}{K} + \frac{(Sc(1 + \alpha_1) - 1)}{2} \right) \sqrt{\frac{Sc \sqrt{Sc^2 + n^2} + Sc^2}{2}} \\
 A_3 &= \frac{Pr^2}{4} - \frac{1}{4} \sqrt{\frac{\sqrt{(Pr^2 + 4R)^2 + Pr^2 n^2} + (Pr^2 + 4R)}{2}} + R
 \end{aligned}$$

$$B_3 = \frac{1}{4} \sqrt{\frac{\sqrt{(\text{Pr}^2 + 4R)^2 + \text{Pr}^2 n^2} - (\text{Pr}^2 + 4R)}{2}}$$

$$A_4 = \left[\frac{-\text{Gr} A A_3}{A_3^2 + B_3^2} - \text{A} J_3 x_3 - \frac{J_3}{K} \right]$$

$$B_4 = \left(\frac{\text{Gr} A B_3}{A_3^2 + B_3^2} \right)$$

$$A_5 = (1 + \alpha_1) \frac{\text{Pr}^2}{2} - \frac{\text{Pr}}{2} + R(1 + \alpha_1) - M - \frac{1}{K} + (\text{Pr}(1 + \alpha_1) - 1) \frac{\sqrt{\text{Pr}^2 + 4R}}{2}$$

$$B_5 = \frac{n}{4}$$

$$A_6 = \text{A} J_4 \text{Sc} - \frac{J_4}{K}, \quad B_6 = \frac{4 \text{Gc} A}{\text{Sc} n}$$

$$A_7 = (1 + \alpha_1) \text{Sc}^2 - \text{Sc} - M - \frac{1}{K}$$

$$A_8 = \text{A} J_5 \beta_1 - \frac{J_5}{K}, \quad B_7 = \frac{8\alpha_1 \text{A} \beta_1 u_0''(0)}{n}$$

$$A_9 = (1 + \alpha_1) \beta_1^2 - \beta_1 - M - \frac{1}{K}, \quad A_{10} = \frac{4 b_4 m_7 x_9^2}{n}$$

$$A_{11} = \alpha_1 \beta_1 \left(\beta_1 + \sqrt{\frac{\beta_1 \sqrt{\beta_1^2 + n^2} + \beta_1^2}{2}} \right)$$

$$B_8 = \alpha_1 \beta_1 \left(\beta_1 + \sqrt{\frac{\beta_1 \sqrt{\beta_1^2 + n^2} - \beta_1^2}{2}} \right)$$

$$A_{12} = \left((1 + \alpha_1) \frac{\beta_1^2}{2} - \frac{\beta_1}{2} - M - \frac{1}{K} + \frac{(\beta_1(1 + \alpha_1) - 1)}{2} \sqrt{\frac{\beta_1 \sqrt{\beta_1^2 + n^2} + \beta_1^2}{2}} \right)$$

$$B_9 = \left((\beta_1(1 + \alpha_1) - 1) \frac{n}{4} + \frac{(\beta_1(1 + \alpha_1) - 1)}{2} \sqrt{\frac{\beta_1 \sqrt{\beta_1^2 + n^2} - \beta_1^2}{2}} \right)$$

$$A_{13} = -\frac{1}{2} \left(\beta_1 + \sqrt{\frac{\beta_1 \sqrt{\beta_1^2 + n^2} + \beta_1^2}{2}} \right), \quad B_{10} = -\frac{1}{2} \sqrt{\frac{\beta_1 \sqrt{\beta_1^2 + n^2} - \beta_1^2}{2}}$$

$$A_{15} = \left\{ \left(\frac{\text{Gr} A x_3 A_1}{A_1^2 + B_1^2} \right) \left(-\frac{\text{Pr}}{2} - \frac{1}{2} \sqrt{\frac{\sqrt{(\text{Pr}^2 + 4R)^2 + \text{Pr}^2 n^2} + (\text{Pr}^2 + 4R)}{2}} \right) \right. \\ \left. - \left(\frac{\text{Gr} A x_3 B_1}{A_1^2 B_1^2} \right) \left(-\frac{1}{2} \sqrt{\frac{\sqrt{(\text{Pr}^2 + 4R)^2 + \text{Pr}^2 n^2} - (\text{Pr}^2 + 4R)}{2}} \right) \right\}$$

$$B_{12} = \left(\frac{Gr A x_3 A_1}{A_1^2 + B_1^2} \right) \left(-\frac{Pr}{2} - \frac{1}{2} \sqrt{\frac{(\Pr^2 + 4R)^2 + \Pr^2 n^2 + (\Pr^2 + 4R)}{2}} \right) + \left(\frac{Gr A x_3 A_1}{A_1^2 + B_1^2} \right) \left(\frac{1}{2} \sqrt{\frac{(\Pr^2 + 4R)^2 + \Pr^2 n^2 - (\Pr^2 + 4R)}{2}} \right)$$

$$A_{16} = \left(\frac{-Gc A A_2}{A_2^2 + B_2^2} \right) \left(-\frac{Sc}{2} - \frac{1}{2} \sqrt{\frac{Sc\sqrt{Sc^2 + n^2} + Sc^2}{2}} \right) + \left(\frac{Gc A B_2}{A_2^2 + B_2^2} \right) \left(\frac{1}{2} \sqrt{\frac{Sc\sqrt{Sc^2 + n^2} - Sc^2}{2}} \right)$$

$$B_{13} = \left(\frac{Gc A B_2}{A_2^2 + B_2^2} \right) \left(-\frac{Sc}{2} - \frac{1}{2} \sqrt{\frac{Sc\sqrt{Sc^2 + n^2} + Sc^2}{2}} \right) + \left(\frac{Gc A A_2}{A_2^2 + B_2^2} \right) \left(\frac{1}{2} \sqrt{\frac{Sc\sqrt{Sc^2 + n^2} - Sc^2}{2}} \right)$$

$$A_{17} = \left[A_1 \left(-\frac{Pr}{2} - \frac{1}{2} \sqrt{\frac{(\Pr^2 + 4r)^2 + \Pr^2 n^2 + (\Pr^2 + 4R)}{2}} + \frac{B_1}{2} \sqrt{\frac{(\Pr^2 + 4R)^2 + \Pr^2 n^2 - (\Pr^2 + 4R)}{2}} \right) \right]$$

$$B_{14} = \left[B_1 \left(-\frac{Pr}{2} - \frac{1}{2} \sqrt{\frac{(\Pr^2 + 4r)^2 + \Pr^2 n^2 + (\Pr^2 + 4R)}{2}} + \frac{A_1}{2} \sqrt{\frac{(\Pr^2 + 4R)^2 + \Pr^2 n^2 - (\Pr^2 + 4R)}{2}} \right) \right]$$

$$A_{18} = \left[\frac{A_2}{2} \left(-Sc - \sqrt{\frac{Sc\sqrt{Sc^2 + n^2} + Sc^2}{2}} + \frac{B_2}{2} \sqrt{\frac{Sc\sqrt{Sc^2 + n^2} - Sc^2}{2}} \right) \right]$$

$$B_{15} = \left[\frac{B_2}{2} \left(-Sc - \sqrt{\frac{Sc\sqrt{Sc^2 + n^2} + Sc^2}{2}} - \frac{A_2}{2} \sqrt{\frac{Sc\sqrt{Sc^2 + n^2} - Sc^2}{2}} \right) \right]$$

$$s_1 = \frac{Gr A x_3 A_{17}}{A_{17}^2 + B_{14}^2}, \quad s_2 = \frac{Gr A x_3 B_{14}}{A_{17}^2 + B_{14}^2}$$

$$d_{46} = s_1 \left(+\frac{h_1}{2} d_{10} - 1 \right) + s_2 \frac{h_1}{2} d_{11}, \quad d_{47} = -s_2 \left(+\frac{h_1}{2} d_{10} - 1 \right) + s_1 \frac{h_1}{2} d_{11}$$

$$s_3 = -\frac{GcAA_{18}}{A_{18}^2 + B_{15}^2}, \quad s_4 = \frac{GcAB_{15}}{A_{18}^2 + B_{15}^2}$$

$$d_{48} = s_3 \left(+\frac{h_1}{2} d_{12} - 1 \right) - s_4 \left(\frac{h_1}{2} d_{13} \right)$$

$$d_{49} = s_4 \left(+\frac{h_1}{2} d_{12} - 1 \right) + s_3 \left(\frac{h_1}{2} d_{13} \right),$$

$$s_7 = \frac{(A_4A_5 + B_4B_5)}{A_5^2 + B_5^2}, \quad s_8 = \frac{(A_4B_5 - A_5B_4)}{A_5^2 + B_5^2}$$

$$d_{52} = s_7 (h_1x_3 - 1), \quad d_{53} = s_8 (h_1x_3 - 1)$$

$$s_9 = \frac{A_6A_7 + B_6B_5}{A_7^2 + B_5^2}, \quad s_{10} = \frac{(A_6B_5 - A_7B_6)}{A_7^2 + B_5^2}$$

$$d_{54} = -s_9 (h_1Sc + 1), \quad d_{55} = s_{10} (h_1Sc + 1)$$

$$s_{11} = \frac{(A_8A_9 - B_7B_5)}{A_9^2 + B_5^2}, \quad s_{12} = \frac{(A_8B_5 + A_9B_7)}{A_9^2 + B_5^2}$$

$$d_{56} = -s_{11} (h_1\beta_1 + 1), \quad d_{57} = s_{12} (h_1\beta_1 + 1)$$

$$d_{58} = \frac{4b_4m_7}{n} (h_1x_9 - 1)$$

$$s_{15} = \frac{-1}{(1 + \alpha_1)} \left(1 + \sqrt{\frac{\left(1 + 4\left(M + \frac{1}{K}\right)(1 + \alpha_1)\right)^2 + n^2(1 + \alpha_1)^2 + \left(1 + 4\left(M + \frac{1}{K}\right)(1 + \alpha_1)\right)}{2}} \right)$$

$$s_{16} = \frac{-1}{(1 + \alpha_1)} \left(1 + \sqrt{\frac{\left(1 + 4\left(M + \frac{1}{K}\right)(1 + \alpha_1)\right)^2 + n^2(1 + \alpha_1)^2 - \left(1 + 4\left(M + \frac{1}{K}\right)(1 + \alpha_1)\right)}{2}} \right)$$

$$d_{60} = 1 - \frac{h_1s_{15}}{2}, \quad d_{61} = \frac{h_1s_{16}}{2}$$

$$d_{62} = (d_{46} + d_{48} + d_{52} + d_{54} + d_{56}), \quad d_{63} = (d_{47} + d_{49} + d_{53} - d_{55} - d_{57} + d_{58})$$

$$d_{64} = \frac{(d_{62}d_{60} + d_{63}d_{61})}{d_{60}^2 + d_{61}^2}, \quad d_{65} = \frac{(d_{60}d_{63} - d_{61}d_{62})}{d_{60}^2 + d_{61}^2}$$

$$l_1 = \frac{s_{15}^2}{4} - \frac{s_{16}^2}{4}, \quad l_2 = s_{15} \left(\frac{s_{16}}{2} \right)$$

$$l_3 = (d_{64}l_1 + d_{65}l_2), \quad l_4 = (d_{65}l_1 - d_{64}l_2)$$

$$l_5 = A_{13}^2 - B_{10}^2$$

$$l_6 = 2A_{13} B_{10}$$

$$l_7 = \left(-\frac{h_1}{2} d_{14} - 1 \right) \quad , \quad l_8 = \frac{h_1 d_{15}}{2}$$

$$d_{14} = \left(\beta_1 + \sqrt{\frac{\beta_1 \sqrt{\beta_1^2 + n^2} + \beta_1^2}{2}} \right) \quad , \quad d_{15} = \sqrt{\frac{\beta_1 \sqrt{\beta_1^2 + n^2} - \beta_1^2}{2}}$$

$$l_9 = \left(1 - \frac{h_1 s_{15}}{2} \right) \quad , \quad l_{10} = \frac{h_1 s_{16}}{2}$$

$$l_{11} = (A_{12} l_9 - B_9 l_{10}) - 2\alpha_1 (l_5 l_9 - l_6 l_{10}) - 2\alpha_1 (l_1 l_7 - l_2 l_8)$$

$$l_{12} = (B_9 l_9 + A_{12} l_{10}) - 2\alpha_1 (l_6 l_9 + l_5 l_{10}) + 2\alpha_1 (l_2 l_7 + l_1 l_8)$$

$$l_{13} = (A_{12} l_9 + B_9 l_{10}) \quad , \quad l_{14} = (B_9 l_9 + A_{12} l_{10})$$

$$d_1 = A_{13} l_{13} + l_{14} B_{10} \quad , \quad d_2 = A_{13} l_{14} - B_{10} l_{13}$$

$$d_3 = A_{13}^2 l_{11} + B_{10}^2 l_{11} \quad , \quad l_{15} = A_{13}^2 l_{12} + B_{10}^2 l_{12}$$

$$l_{22} = \frac{d_1 d_3 + d_2 l_{15}}{d_3^2 + l_{15}^2} \quad , \quad l_{23} = \frac{d_2 d_3 - d_1 l_{15}}{d_3^2 + l_{15}^2}$$

$$d_4 = \left\{ -l_3 - A_{15} - A_{16} - \frac{x_3^2 (A_4 A_5 + B_4 B_5)}{A_5^2 + B_5^2} - \frac{Sc^2 (A_6 A_7 + B_6 B_5)}{A_7^2 + B_5^2} - \frac{\beta_1^2 (A_8 A_9 - B_7 B_5)}{A_9^2 + B_5^2} \right\}$$

$$d_5 = \left\{ -l_4 - B_{12} - B_{13} - \frac{x_3^2 (A_4 B_5 - A_5 B_4)}{A_5^2 + B_5^2} - \frac{Sc^2 (A_6 B_5 - A_7 B_6)}{A_7^2 + B_5^2} - \frac{\beta_1^2 (A_8 B_5 + A_9 B_7)}{A_9^2 + B_5^2} - A_{10} + \frac{4u_0''(0)\beta_1}{n} \right\}$$

$$d_6 = (l_{22} d_4 - l_{23} d_5) \quad , \quad d_7 = (l_{23} d_4 + l_{22} d_5)$$

$$d_8 = -\beta_1 - \sqrt{\frac{\beta_1 \sqrt{\beta_1^2 + n^2} + \beta_1^2}{2}} \quad , \quad d_9 = -\sqrt{\frac{\beta_1 \sqrt{\beta_1^2 + n^2} - \beta_1^2}{2}}$$

$$L_3 = \frac{d_9 y}{2} \quad , \quad C_1 = \cos L_3 \quad , \quad C_2 = \sin L_3$$

$$L_2 = -\frac{s_{16} y}{2} \quad , \quad C_7 = \cos L_2 \quad , \quad C_8 = \sin L_2$$

$$d_{66} = e^{s_{15} y/2} (d_{64} C_7 - d_{65} C_8) \quad , \quad d_{67} = e^{s_{15} y/2} (d_{65} C_7 + d_{64} C_8)$$

$$d_{10} = \left(-Pr - \sqrt{\frac{(\text{Pr}^2 + 4R)^2 + \text{Pr}^2 n^2 + (\text{Pr}^2 + 4R)}{2}} \right)$$

$$d_{11} = \left(-\sqrt{\frac{(\text{Pr}^2 + 4R)^2 + \text{Pr}^2 n^2 - (\text{Pr}^2 + 4R)}{2}} \right)$$

$$L_4 = \frac{d_{11}y}{2}, \quad C_3 = \cos L_4, \quad C_4 = \sin L_4$$

$$d_{16} = e^{d_{10}y/2} (s_1 C_3 + s_2 C_4)$$

$$d_{17} = e^{d_{10}y/2} (s_1 C_4 - s_2 C_3)$$

$$d_{12} = -Sc - \sqrt{\frac{Sc\sqrt{Sc^2 + n^2} + Sc^2}{2}}, \quad d_{13} = -\sqrt{\frac{Sc\sqrt{Sc^2 + n^2} - Sc^2}{2}}$$

$$L_5 = \frac{d_{13}y}{2}, \quad C_5 = \cos L_5, \quad C_6 = \sin L_5$$

$$d_{18} = e^{d_{12}y/2} (s_3 C_5 - s_4 C_6), \quad d_{19} = e^{d_{12}y/2} (s_4 C_5 + s_3 C_6)$$

$$l_{24} = \left(-\frac{d_6 d_8}{2} + \frac{d_7 d_9}{2} \right), \quad l_{25} = \left(-\frac{d_6 d_9}{2} - \frac{d_7 d_8}{2} \right)$$

$$l_{26} = l_{24} A_{12} + l_{25} B_9, \quad l_{27} = (A_{12} l_{25} - B_9 l_{24})$$

$$l_{28} = \frac{l_{26}}{A_{12}^2 + B_9^2}, \quad l_{29} = \frac{l_{27}}{A_{12}^2 + B_9^2}$$

$$d_{20} = e^{d_8 y/2} (l_{28} C_1 - l_{29} C_2), \quad d_{21} = e^{d_8 y/2} (l_{29} C_1 + l_{28} C_2)$$

$$d_{22} = s_7 e^{x_3 y}, \quad d_{23} = s_8 e^{x_3 y}$$

$$d_{24} = s_9 \bar{e}^{Scy}, \quad d_{25} = s_{10} \bar{e}^{Scy}$$

$$d_{26} = s_{11} \bar{e}^{\beta_1 y}, \quad d_{27} = s_{12} \bar{e}^{\beta_1 y}$$

$$d_{28} = \frac{4b_4 m_7}{n} e^{x_9 y}$$

$$d_{30} = -\left(\frac{Pr n}{4} \right) \sqrt{\frac{\sqrt{(Pr^2 + 4r)^2 + Pr^2 n^2} - (Pr^2 + 4R)}{2}}$$

$$d_{31} = \left(\frac{Pr n}{4} \right) \left(\frac{Pr}{2} + \frac{1}{2} \sqrt{\frac{\sqrt{(Pr^2 + 4R)^2 + Pr^2 n^2} + (Pr^2 + 4R)}{2}} \right)$$

$$l_{30} = \frac{-Ax_3 d_{30}}{d_{30}^2 + d_{31}^2}, \quad l_{31} = \frac{Ax_3 d_{31}}{d_{30}^2 + d_{31}^2}$$

$$d_{34} = e^{d_{10}y/2} \{l_{30} C_3 - l_{31} C_4\}$$

$$d_{35} = e^{d_{10}y/2} \{l_{31} C_3 + l_{30} C_4\}$$

$$l_{32} = \frac{AA_3}{A_3^2 + B_3^2}, \quad l_{33} = \frac{AB_3}{A_3^2 + B_3^2}$$

$$d_{36} = l_{32} e^{x_3 y}, \quad d_{37} = l_{33} e^{x_3 y}$$

$$d_{38} = -\left(\frac{Scn}{4} \right) \left(\frac{1}{2} \sqrt{\frac{Sc\sqrt{Sc^2 + n^2} - Sc^2}{2}} \right)$$

$$d_{39} = \left(\frac{Scn}{4} \right) \left(\frac{Sc}{2} + \frac{1}{2} \sqrt{\frac{Sc\sqrt{Sc^2 + n^2} + Sc^2}{2}} \right)$$

$$l_{34} = \frac{ASc d_{38}}{d_{30}^2 + d_{39}^2}, \quad l_{35} = \frac{ASc d_{39}}{d_{30}^2 + d_{39}^2}$$

$$C_5 = \cos \frac{d_{13}y}{2}, \quad C_6 = \sin \frac{d_{13}y}{2}$$

$$d_{42} = e^{d_{12}y/2} (l_{34}C_5 + l_{35}C_6)$$

$$d_{43} = e^{d_{12}y/2} (l_{34}C_6 - l_{35}C_5)$$

$$d_{44} = \frac{4A}{Scn} e^{-Scy}$$

$$t_1 = \left(d_{64} \frac{s_{15}}{2} + d_{65} \frac{s_{16}}{2} \right), \quad t_2 = \left(d_{65} \frac{s_{15}}{2} - d_{64} \frac{s_{16}}{2} \right)$$

$$t_3 = \left\{ s_1 \left(-\frac{Pr}{2} - \frac{1}{2} \sqrt{\frac{(\Pr^2 + 4R)^2 + \Pr^2 n^2 + (\Pr^2 + 4R)}{2}} \right) \right. \\ \left. -s_2 \left(\frac{1}{2} \sqrt{\frac{(\Pr^2 + 4R)^2 + \Pr^2 n^2 - (\Pr^2 + 4R)}{2}} \right) \right\}$$

$$t_4 = \left\{ -s_2 \left(-\frac{Pr}{2} - \frac{1}{2} \sqrt{\frac{(\Pr^2 + 4R)^2 + \Pr^2 n^2 + (\Pr^2 + 4R)}{2}} \right) \right. \\ \left. -s_1 \left(\frac{1}{2} \sqrt{\frac{(\Pr^2 + 4R)^2 + \Pr^2 n^2 - (\Pr^2 + 4R)}{2}} \right) \right\}$$

$$t_5 = \left\{ s_3 \left(-\frac{Sc}{2} - \frac{1}{2} \sqrt{\frac{Sc\sqrt{Sc^2 + n^2} + Sc^2}{2}} \right) + s_4 \left(\frac{1}{2} \sqrt{\frac{Sc\sqrt{Sc^2 + n^2} - Sc^2}{2}} \right) \right\}$$

$$t_6 = \left\{ s_4 \left(-\frac{Sc}{2} - \frac{1}{2} \sqrt{\frac{Sc\sqrt{Sc^2 + n^2} + Sc^2}{2}} \right) - s_3 \left(\frac{1}{2} \sqrt{\frac{Sc\sqrt{Sc^2 + n^2} - Sc^2}{2}} \right) \right\}$$

$$t_7 = \left\{ l_{28} \left(-\frac{\beta_1}{2} - \frac{1}{2} \sqrt{\frac{\beta_1\sqrt{\beta_1^2 + n^2} + \beta_1^2}{2}} \right) + l_{29} \left(\frac{1}{2} \sqrt{\frac{\beta_1\sqrt{\beta_1^2 + n^2} - \beta_1^2}{2}} \right) \right\}$$

$$t_8 = \left\{ l_{29} \left(-\frac{\beta_1}{2} - \frac{1}{2} \sqrt{\frac{\beta_1 \sqrt{\beta_1^2 + n^2} + \beta_1^2}{2}} \right) - l_{28} \left(\frac{1}{2} \sqrt{\frac{\beta_1 \sqrt{\beta_1^2 + n^2} - \beta_1^2}{2}} \right) \right\}$$

$$t_9 = s_7 x_3 \quad , \quad t_{10} = s_8 x_3$$

$$t_{11} = s_9 Sc \quad , \quad t_{12} = s_{10} Sc$$

$$t_{13} = s_{11} \beta_1 \quad , \quad t_{14} = s_{12} \beta_1$$

$$t_{15} = d_{58} x_9 \quad , \quad t_{16} = (t_1 + t_3 + t_5 + t_7 + t_9 - t_{11} - t_{13})$$

$$t_{17} = (t_2 + t_4 + t_6 + t_8 + t_{10} + t_{12} - t_{14} + t_{15})$$

$$t_{18} = (l_{30} + l_{32}) \quad , \quad t_{19} = (l_{31} + l_{33})$$

$$t_{20} = m_3 + \epsilon (t_{18} \cos nt - t_{19} \sin nt)$$

$$t_{21} = (t_{18} \sin nt + t_{19} \cos nt)$$

VI. Result and Discussion

In order to get a physical insight of the problem, calculations have been made for velocity (u), angular velocity (ω), temperature (θ), concentration (C), skin friction (τ_w) and Nusselt number (Nu) for different values of K (permeability parameter), M (magnetic parameter), Gr (thermal Grashof number), Gc (mass Grashof number), h_1 (velocity slip parameter), R (radiation parameter) etc. Observations are made for both the basic fluids, air ($Pr = 0.71$, $Sc = 0.22$) and water ($Pr = 7$, $Sc = 0.61$) fixing $n = 3$, $nt = \pi/2$, $\epsilon = 0.01$.

The velocity profiles are plotted against y in figures 1, 2 and 3 taking $B = 0.5$. We observe that for both the basic fluids air and water, velocity increases with increase in K , Gr , Gc and β_1 where as on increasing M and β_1 velocity drops. On negative of radiation i.e, when absorption takes place, it tends the velocity to rise. It is note worthy, that on decreasing the velocity slip parameter (h_1), the velocity drops at the plate but rises as we move away from the plate. On the other hand when we decrease A , velocity increases at the plate and drops later. For the case of free flow with no slip at the boundary i.e ($K = \infty$, $h_1 = 0$) as compared to ($K \neq \infty$, $h_1 \neq 0$) velocity drops at the plate but increases as soon as we move away from the plate. The effect of K can be understood physically as, when K increases, that is, the medium is more porous, so more fluid can flow through hence increasing the velocity of the fluid. On the whole we observe that velocity is less for water as compared to the case for air, this is because air is lighter than water.

In figure 4, angular velocity profiles are plotted against y fixing $A = B = 0.5$, $R = 0.05$, $Gr = 0.3$, $Gc = 0.1$ and $M = 2$. From the figure it is observed that for both the basic fluids air and water, on increasing K and β_1 , angular velocity decreases but, angular velocity increases with increase in h_1 and β_1 , for the case of free flow, with no slip at the boundary ($K = \infty$, $h_1 = 0$), angular velocity drops, as compared to the case of ($K \neq \infty$, $h_1 \neq 0$). In general we notice that angular velocity is lower for air as compared to water.

Temperature profiles are plotted against y , in figure 5. We observe for both the basic fluids air and water increase in radiation, decreases the temperature and negative of radiation i.e absorption tends the temperature to rise as when absorption of heat takes place the temperature of the fluid will rise. But the effect of A is different for the case of air as compared to water, since for air, when suction velocity is constant ($A = 0$), the temperature rises, where as, for the same temperature drops for water.

In figure 6, concentration profiles are plotted against y , for different values of Schmidt number (Sc), chosen in such a way that they represent the diffusing chemical species of most common interest in air (for example, the values of Schmidt number for H_2 , H_2O , NH_3 and propyl benzene in air is 0.22, 0.61, 0.78 and 2.62 respectively). We clearly observe that as Sc increases, concentration decreases. We also notice that for constant suction velocity ($A = 0$), concentration increases slightly for both the basic fluids air and water. Moreover, for every value of A , concentration is much higher for air as compared to water.

An important physical parameter skin friction is plotted against K , in figures 7 and 8. It is noteworthy from the figures that for both the basic fluids air and water, increase in Gr , Gc and β_1 , increases the skin friction, where as decrease in M , h_1 and β_1 , increases the skin friction. The effect of velocity slip parameter (h_1), can be seen physically as when h_1 decreases i.e slip at the boundary is decreased, this will lead to increase in

skin friction, which exactly happens in our case for both the basic fluids. From both the figures we observe that skin friction is higher for air as compared to water.

Another important physical parameter, Nusselt number is plotted against A in figure 9. It is observed that for both the basic fluids air and water, increase in radiation increases the Nusselt number and when absorption (negative of radiation) takes place Nusselt number drops. This is because when radiation increases, the rate of heat transfer is high thus Nusselt number increases. It is noteworthy that Nusselt number is higher for water as compared to air.

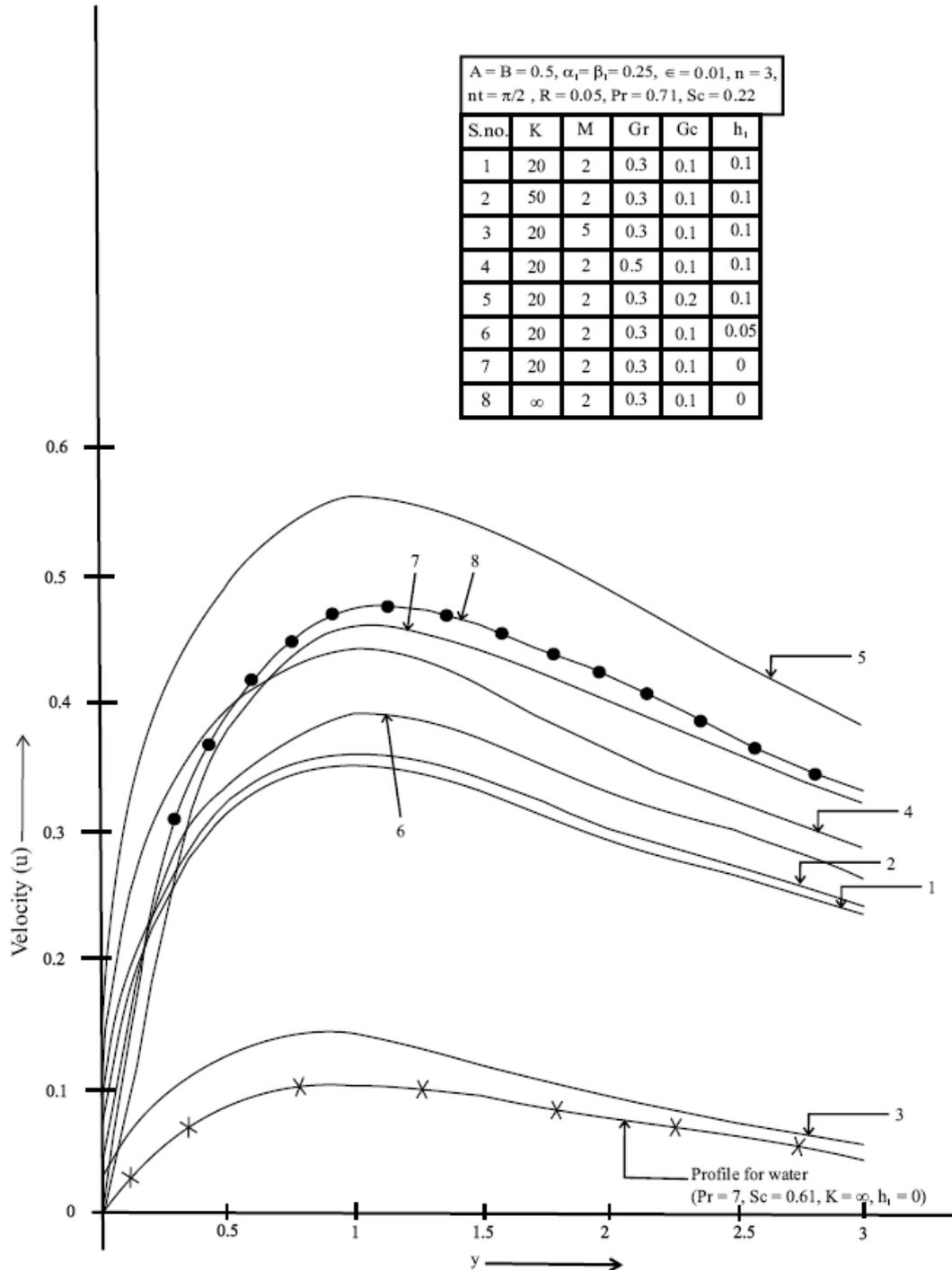


Figure 1: Velocity profiles for air plotted against y for different values of K, M, Gr, Gc and h_1 .

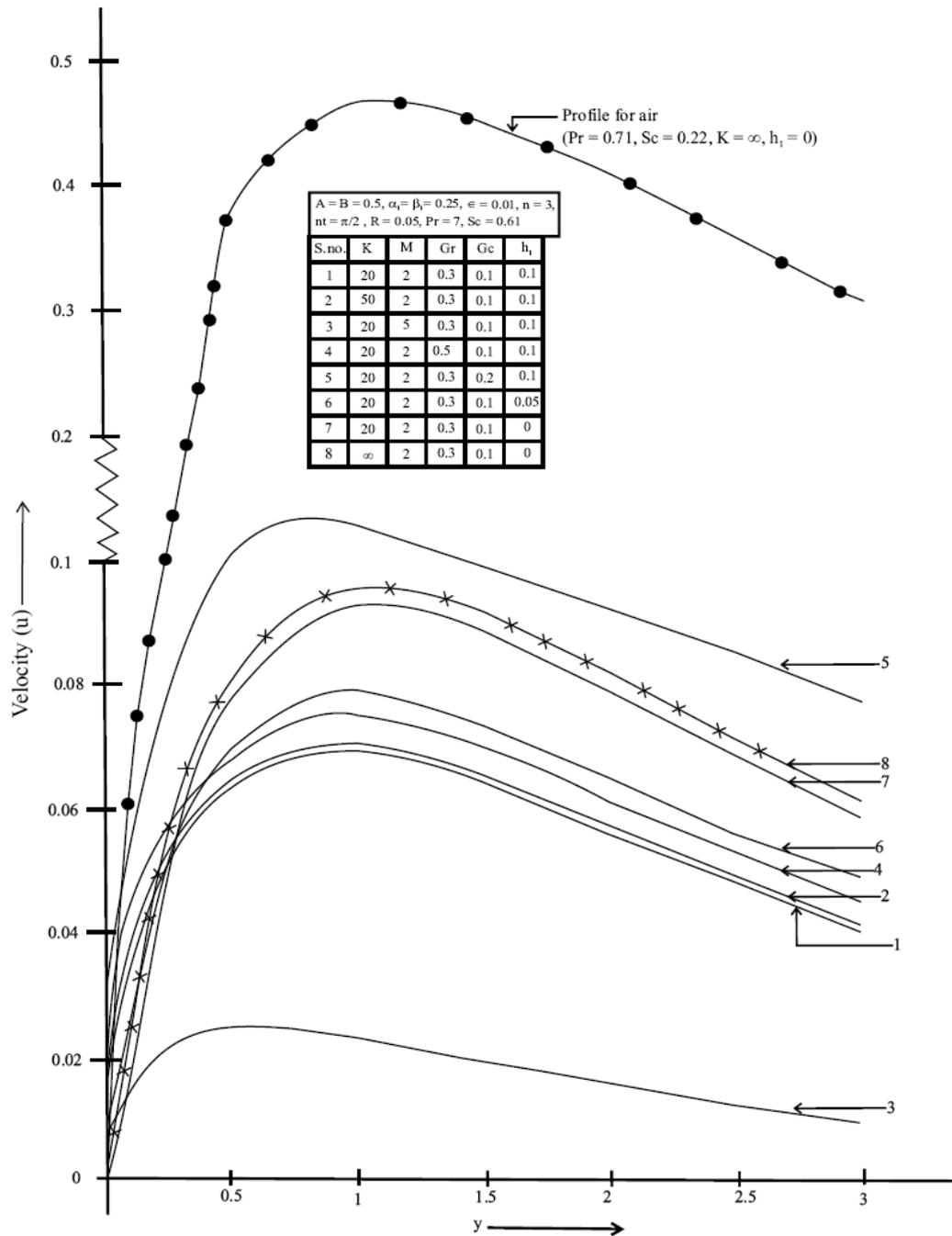


Figure 2: Velocity profiles for water plotted against y for different values of K, M, Gr, Gc and h₁.

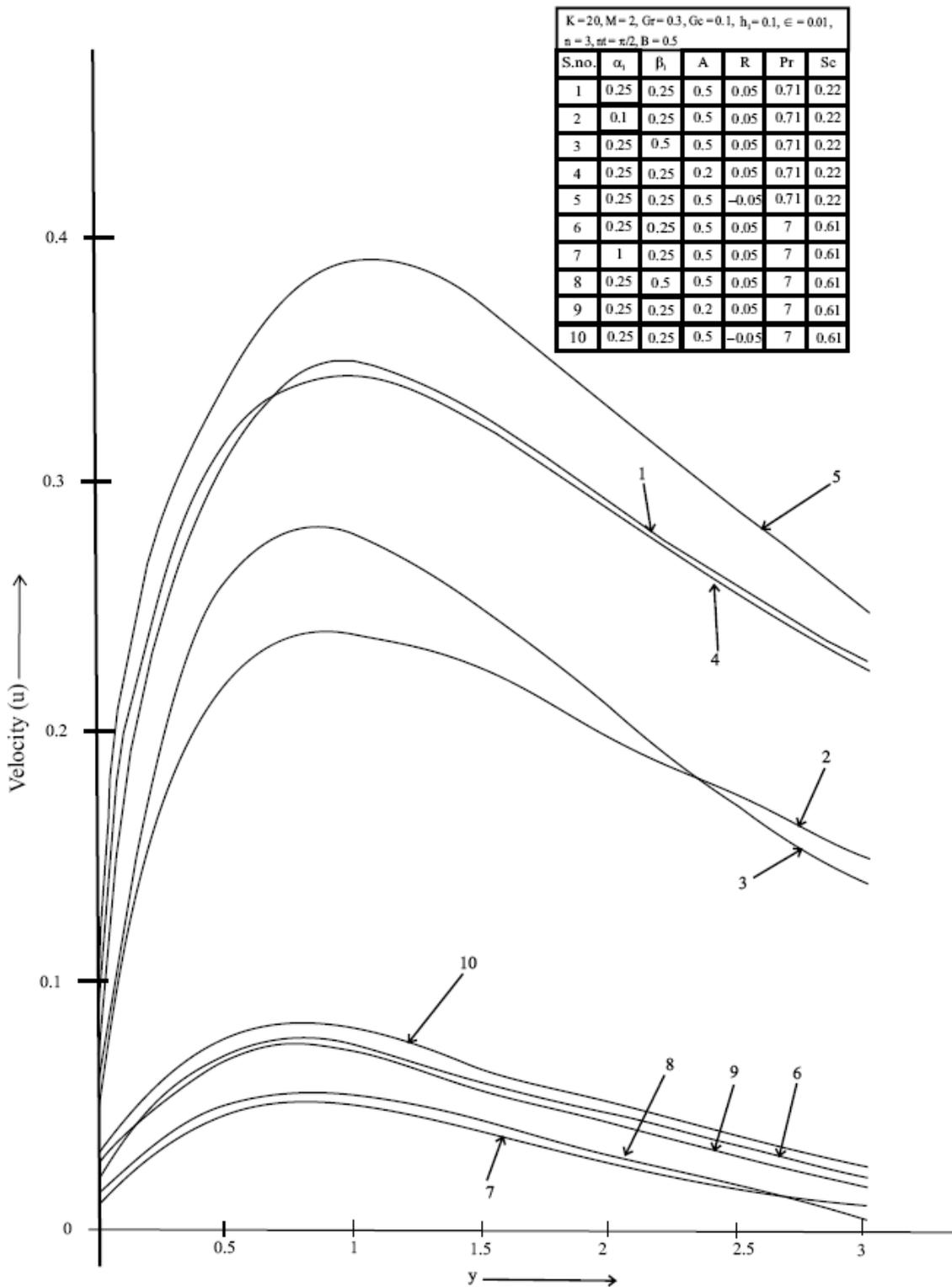


Figure 3: Velocity profiles plotted against y for different values of α_1 , β_1 , A, R, Pr and Sc.

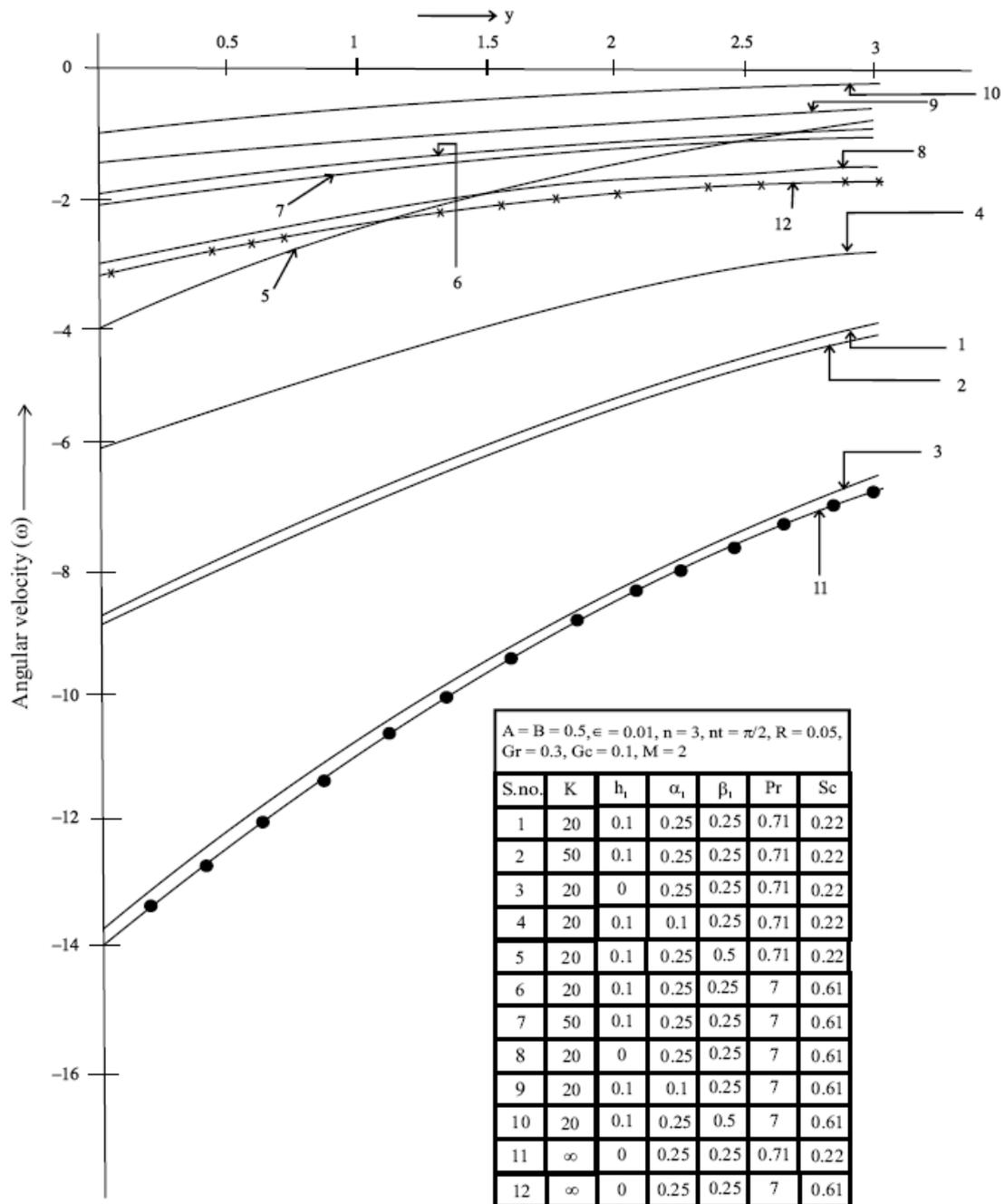


Figure 4: Angular velocity profiles plotted against y for different values of $K, h_1, \alpha_1, \beta_1, Pr$ and Sc .

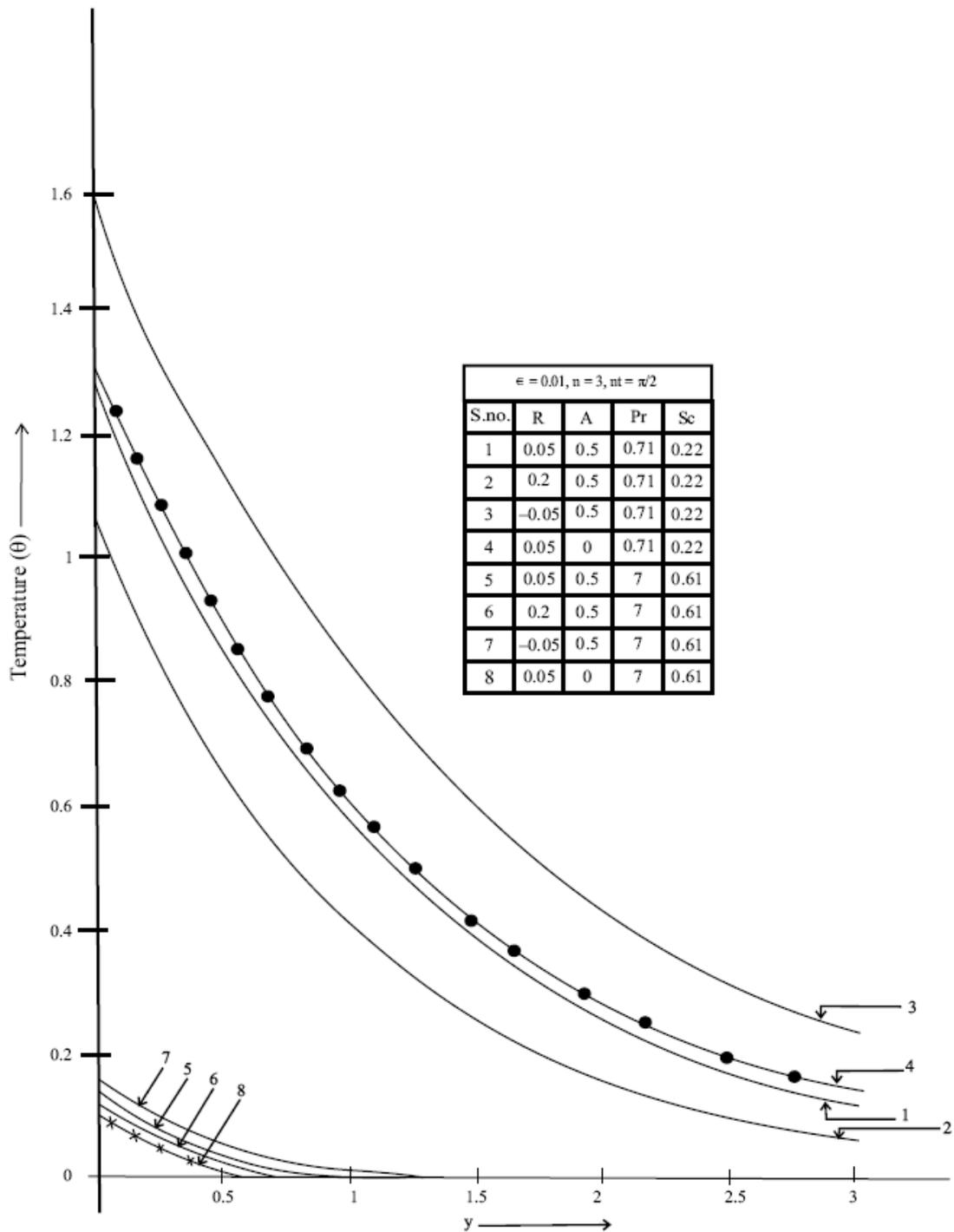


Figure 5: Temperature profiles plotted against y for different values of R, A, Pr and Sc.

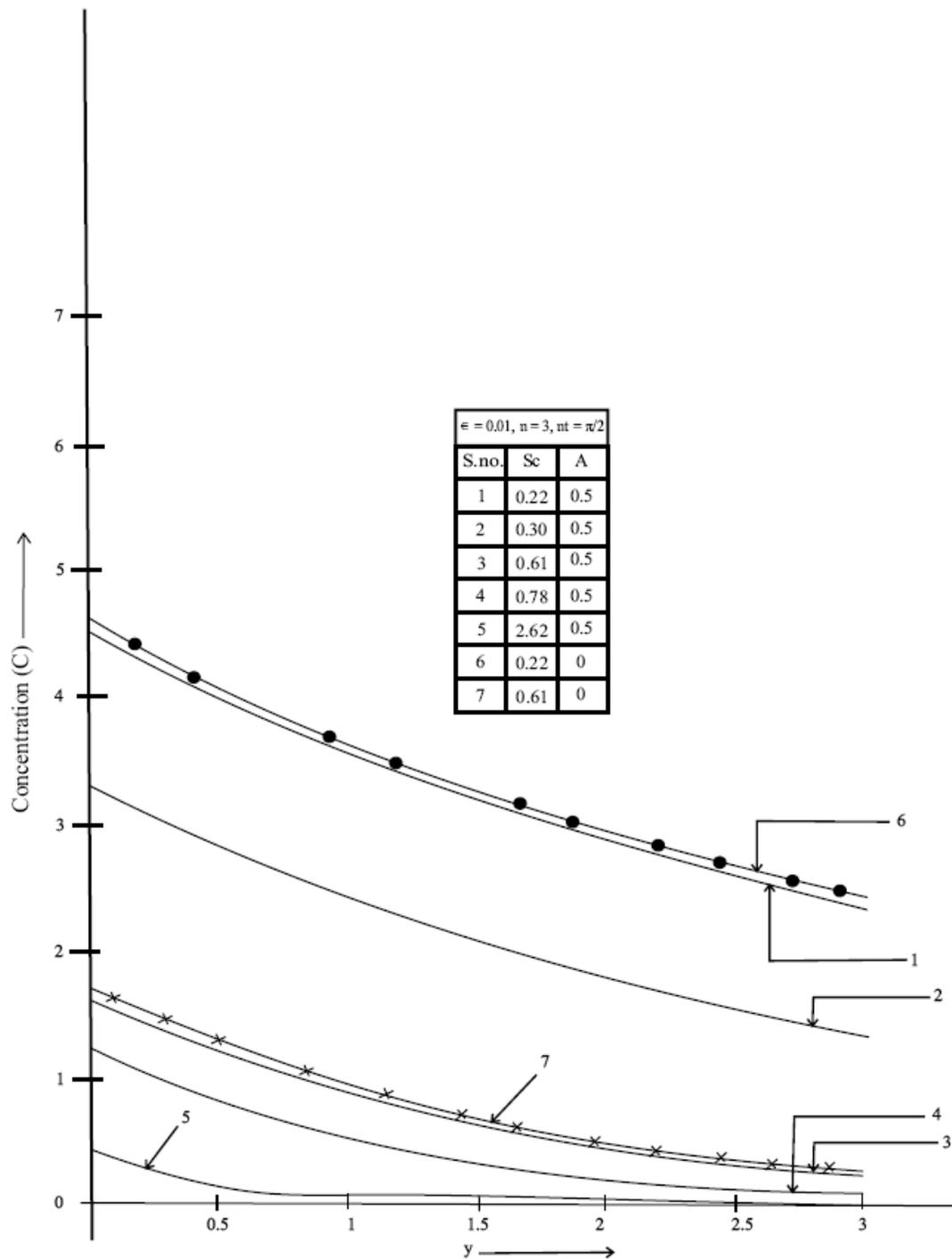


Figure 6: Concentration profiles plotted against y for different values of Sc and A .

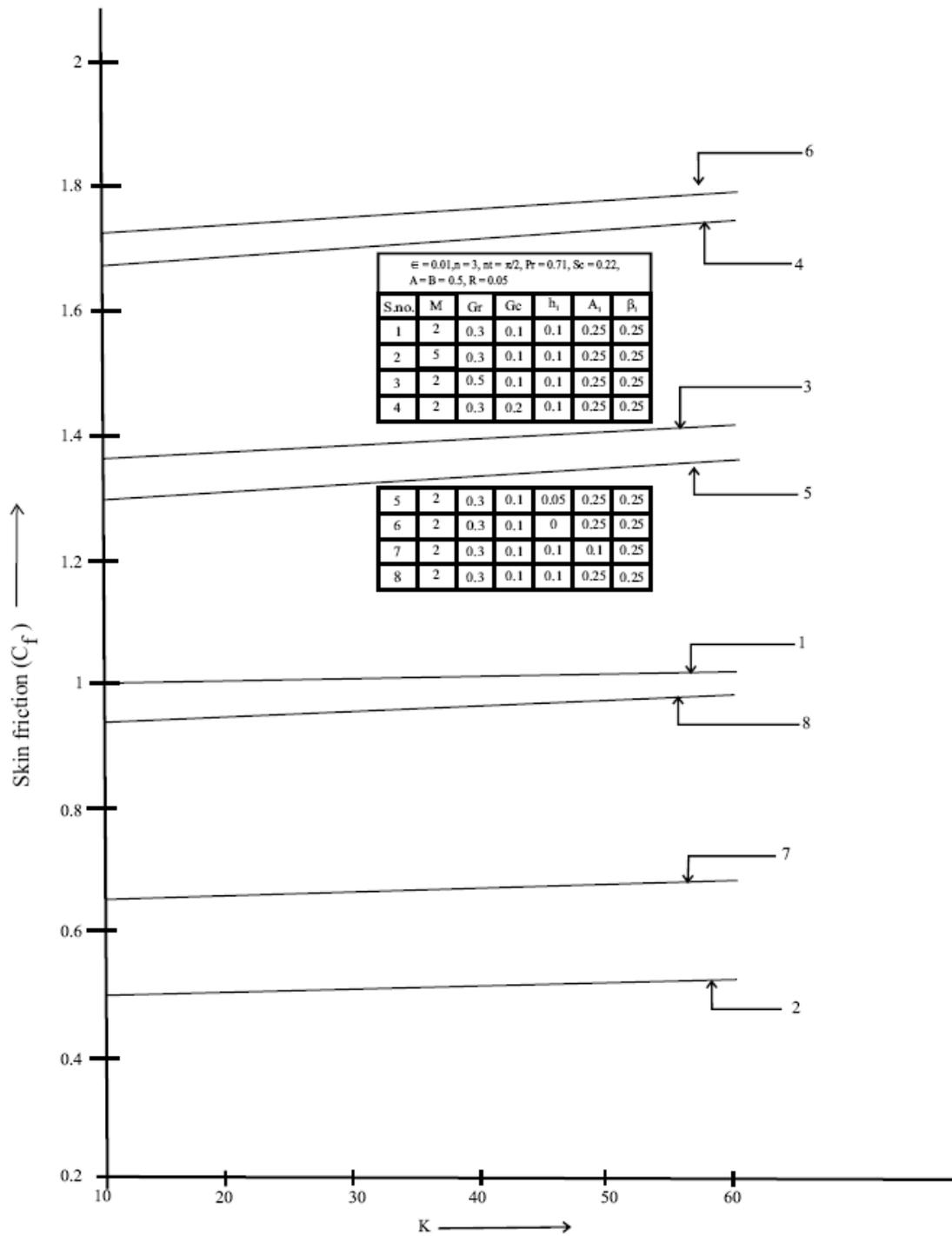


Figure 7: Skin friction for air plotted against K for different values of M, Gr, Gc, h_1 , α_1 , β_1 .

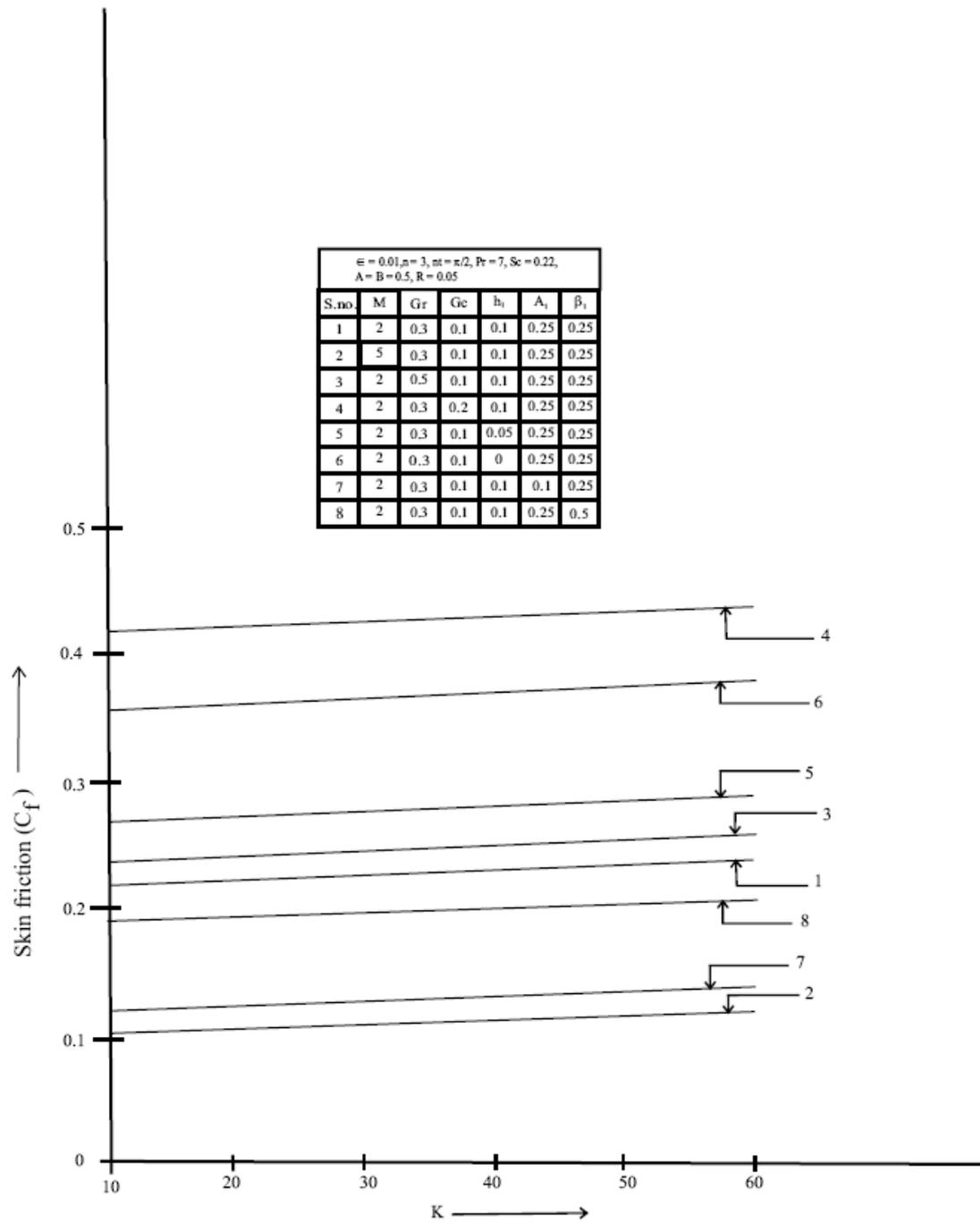


Figure 8: Skin friction for water plotted against K for different values of M, Gr, Gc, h_1 , α_1 , β_1 .

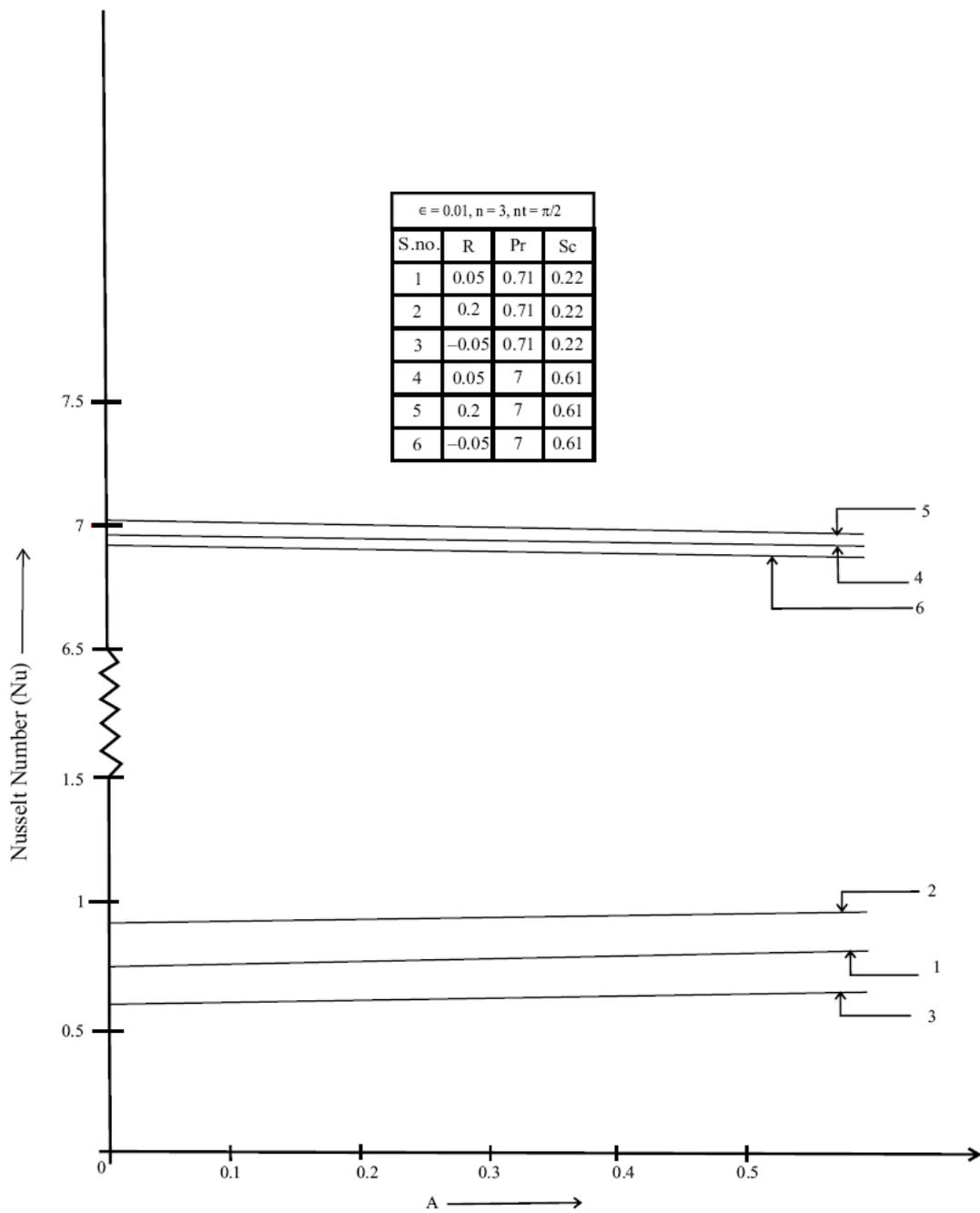


Figure 9: Nusselt number plotted against A for different values of R, Pr and Sc.

1. References

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